# QUANTUM CHEMICAL B3LYP/cc-pvqz COMPUTATION OF GROUND-STATE STRUCTURES AND PROPERTIES OF SMALL MOLECULES WITH ATOMS OF $\mathrm{Z} \leq 18$ (HYDROGEN TO ARGON) 

# (IUPAC Technical Report) 

Prepared for publication by RUDOLF JANOSCHEK<br>Institut für Theoretische Chemie, Karl-Franzens-Universität Graz, Strassoldogasse 10 A-8010 Graz, Austria


#### Abstract

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# Quantum chemical B3LYP/cc-pvqz computation of ground-state structures and properties of small molecules with atoms of $Z \leq 18$ (hydrogen to argon) 

(IUPAC Technical Report)


#### Abstract

Since density functional theory (DFT) achieved a remarkable breakthrough in computational chemistry, the important general question "How reliable are quantum chemical calculations for spectroscopic properties?" should be answered anew. In this project, the most successful density functionals, namely the Becke B3LYP functionals, and the correlation-consistent polarized valence quadruple zeta basis sets (cc-pvqz) are applied to small molecules. In particular, the complete set of experimentally known diatomic molecules formed by the atoms H to Ar (these are 214 species) is uniformly calculated, and calculated spectroscopic properties are compared with experimental ones. Computationally demanding molecules, such as open-shell systems, anions, or noble gas compounds, are included in this study. Investigated spectroscopic properties are spectroscopic ground state, equilibrium internuclear distance, harmonic vibrational wavenumber, anharmonicity, vibrational absolute absorption intensity, electric dipole moment, ionization potential, and dissociation energy. The same computational method has also been applied to the ground-state geometries of 56 polyatomic molecules up to the size of benzene. Special sections are dedicated to nuclear magnetic resonance (NMR) chemical shifts and isotropic hyperfine coupling constants. Each set of systems for a chosen property is statistically analyzed, and the above important question "How reliable...?" is mathematically answered by the mean absolute deviation between calculated and experimental data, as well as by the worst agreement. In addition to presentation of numerous quantum chemically calculated spectroscopic properties, a corresponding updated list of references for experimentally determined properties is presented.


## CONTENTS

1. INTRODUCTION
2. MOLECULAR PROPERTIES, SYMBOLS, AND UNITS
3. COMPUTATIONAL PROCEDURES
4. BASIS SETS
5. STATISTICS OF DEVIATIONS BETWEEN CALCULATED AND EXPERIMENTAL PROPERTIES OF DIATOMIC MOLECULES
6. TABLE OF PROPERTIES OF DIATOMIC MOLECULES
7. ISOTROPIC HYPERFINE COUPLING CONSTANTS OF POLYATOMIC MOLECULES
8. STRUCTURES OF POLYATOMIC GAS-PHASE MOLECULES
9. ANHARMONICITIES OF DIATOMIC MOLECULES
10. NMR CHEMICAL SHIFTS OF POLYATOMIC MOLECULES
11. CONCLUDING REMARKS ON CALCULATED VERSUS EXPERIMENTAL PROPERTIES AND COMPUTATIONALLY DEMANDING MOLECULES
12. REFERENCES

## 1. INTRODUCTION

This project is concerned with experimentally known small molecules, composed of the atoms H up to Ar. Special attention is paid to diatomic molecules where numerous experimentally determined properties are available for comparison with calculated data. The preselected set of 18 atoms gives rise in principle to 171 distinct neutral diatomic molecules. Consideration of the corresponding cations and anions, ignoring any possible excited states, would yield a grand total of some 500 basic molecular systems. Most of these, however, are still hypothetical, since only a subset, which comprises 214 systems out of these basic diatomics, has been observed hitherto. The entire subset of these species for which there is experimental evidence has been calculated uniformly at the B3LYP/cc-pvqz level of density functional theory, and the evaluated properties are compared with the corresponding available experimental values. The presentation of the results is similar to that in the book of Huber and Herzberg [1].

This project is intended to serve a threefold aim. On one hand, the predictive power of recently developed quantum chemical methods will be demonstrated, using the experimental data as a gauge. On the other hand, it will be shown that the calculated properties are very useful to complete our knowledge of these diatomic systems, for which there exists merely inchoate experimental information. Finally, carefully updated sets of experimentally detected molecular properties are presented, irrespective of the theoretical interests of the user. Most importantly, statistics are presented for the accuracy of the respective calculated molecular properties without the omission of computationally difficult molecules.

## 2. MOLECULAR PROPERTIES, SYMBOLS, AND UNITS

The set of 214 diatomic molecules splits into 76 closed-shell and 138 open-shell molecules. The set of 138 open-shell systems splits into 92 radicals with a doublet ground state and 46 radicals with higher spin multiplicities (triplet, quartet, quintet) in their ground states. From the set of 214 diatomic molecules, 102 are neutral species, 73 are cations, and 39 are anions. From 73 cations, 50 are radical cations; from 39 anions, 18 are radical anions. For 20 diatomic systems, low-lying electronic states were calculated. A set of 21 experimentally known noble gas ( $\mathrm{He}, \mathrm{Ne}, \mathrm{Ar}$ ) diatomic molecules were calculated. This situation represents a challenge for experimental as well as computational methods. Electronic states and their symmetry symbols, including the spin multiplicity, are important non-numerical features of molecules. For molecular systems like $\mathrm{CO}, \mathrm{N}_{2}, \mathrm{O}_{2}$, the ground states of which have been unequivocally assigned, experimental as well as computed data are restricted to this single electronic state. In ambiguous cases, where low-lying electronic states are expected, several conceivable states have been investigated. This procedure provides not only the identification of experimentally unknown ground states $\left(\mathrm{AlO}^{+}\right)$, but also confirms $(\mathrm{LiO})$ or rebuts $\left(\mathrm{C}_{2}{ }^{+}\right)$previous tentative assignments of the respective lowest-lying states. It should be noted, however, that the density functional approach fails to predict very weakly bound diatomic van der Waals-complexes such as $\mathrm{He}_{2}$ or $\mathrm{Ne}_{2}$, which require a thorough description of dispersion forces. A few exceptional molecules like $\mathrm{B}_{2}, \mathrm{C}_{2}$, and MgO need multi-configuration wave functions to establish the correct ground states. For dubious cases, ab initio ACPF (averaged coupled pair functional) calculations are recommended for a correct sequence of closely spaced ( 0.1 eV ) electronic states. However, improved computational methods, for example, the widely used coupled cluster $\operatorname{CCSD}(\mathrm{T})$ method or the multireference ACPF approach, combined with extended basis sets, are much more time-consuming (for diatomic molecules a factor of 10 compared with the B3LYP method). The advantage of the B3LYP procedure, chosen for this project, is its recognized success as well as its ability to extend the range of application to systems up to 10 heavy atoms and to investigate hundreds of systems in reasonable time.

Relative term values $T_{\mathrm{e}}$, referring to the ground state $\left(T_{\mathrm{e}}=0\right)$, are the differences of total electronic minimum energies at the equilibrium geometries divided by the Planck constant and the speed of light and are given in $\mathrm{cm}^{-1}$. The equilibrium internuclear distance $r_{\mathrm{e}}$, the harmonic vibrational wave number $\omega_{\mathrm{e}}$, the vibrational absolute absorption intensity $A$, and the electric dipole moment $\mu$ are given
in units of ångström $(\AA), \mathrm{cm}^{-1}, \mathrm{~km} \mathrm{~mol}^{-1}$, and debye (D), respectively (see the attached glossary). The sign of the dipole moment is positive for $\mathrm{A}^{-} \mathrm{B}^{+}$and negative for $\mathrm{A}^{+} \mathrm{B}^{-}$. Calculated dipole moments carry either a plus or a minus sign, or are zero by symmetry. For most of the experimental dipole moments no signs are available. Few exceptions are CN, CO, ClF, and SiS. All calculated dipole moments are $\mu_{\mathrm{e}}$ values. The origin of the dipole moment of ionic systems is chosen as the center of mass. Adiabatic ionization energies $\left(E_{\mathrm{i}}\right)$ are corrected for zero-point vibrations and are given in eV . Ionization energies have been calculated only if there exist an experimental $E_{\mathrm{i}}$ and/or the ionized species has been detected experimentally. The dissociation energy $D_{0}(\mathrm{eV})$ of a molecule refers to the separated atoms in their respective ground states, and is corrected for zero-point energy. The separated atoms for ionic systems are indicated in parentheses. For $\mathrm{AB}^{+}$the lowest dissociation energy $D_{0}\left(\mathrm{~A}+\mathrm{B}^{+}\right)$is adopted if $E_{\mathrm{i}}(\mathrm{A})>$ $E_{\mathrm{i}}(\mathrm{B})$. Therefore, Table 1 with atomic energies and ionization energies is attached to facilitate the identification of the pertinent asymptotes in these cases. Experimental atomic data are taken from the JANAF or Moore's tables [2], or from Bowers [3].

The information about diatomic molecules is summarized in Tables 2 and 3. Separate tables (Tables $7 \mathrm{a}-\mathrm{h}$ ) are given for comparison of experimental and calculated NMR chemical shifts $\delta(\mathrm{ppm})$ for the most abundant nuclei of 50 diatomic and small polyatomic molecules. Table 6 compares calculated and experimental anharmonicities, $\omega_{\mathrm{e}} x_{\mathrm{e}}\left(\mathrm{cm}^{-1}\right)$, for 32 diatomic molecules. Experimental and calculated isotropic hyperfine coupling constants, i.e., Fermi contact parameters, $A_{\text {iso }}(\mathrm{G})$ are compared in Table 4 for 12 diatomic and 8 polyatomic molecules of size up to the benzyl radical. And in Table 5 the calculated structures of 53 polyatomic molecules up to the size of $\mathrm{SF}_{6}$ are compared with experimental structures from MW (microwave spectroscopy) or ED (electron diffraction).

## Glossary

| Symbol | Description | Unit |
| :--- | :--- | :--- |
| $E$ | Total electronic ground-state energy | $E_{\mathrm{h}}$ (hartree) |
| State | Irreducible representation of the symmetry group |  |
| $T_{\mathrm{e}}$ | Relative electronic energy referred to the ground state | $\mathrm{cm}^{-1}$ |
| $r_{\mathrm{e}}$ | Equilibrium internuclear distance | $\AA($ ångström $)$ |
| $\omega_{e}$ | Harmonic vibrational wavenumber | $\mathrm{cm}^{-1}$ |
| $\omega_{\mathrm{e}} x_{\mathrm{e}}$ | Anharmonicity of vibrations | $\mathrm{cm}^{-1}$ |
| $A$ | Vibrational absolute absorption intensity | $\mathrm{km} \mathrm{mol}^{-1}$ |
| $\mu$ | Electric dipole moment | $\mathrm{D}(\mathrm{debye})$ |
| $E_{\mathrm{i}}$ | Adiabatic, zero-point energy corrected ionization energy | eV |
| $D_{0}$ | Dissociation energy, corrected for zero-point energy | eV |
| $\delta$ | NMR (nuclear magnetic resonance) chemical shift | ppm |
| $\sigma$ | NMR absolute shielding | ppm |
| $A_{\text {iso }}$ | Isotropic hyperfine coupling constant | G (gauss) |

## Conversion of units

Length
Energy
Dipole moment (electric)
Absorption intensity
Isotropic hyperfine coupling constant


From Quantities, Units, and Symbols in Physical Chemistry, $2^{\text {nd }}$ ed., I. Mills, T. Cvitaš, K. Homann, N. Kallay and K. Kuchitsu, Blackwell Scientific Publications, Oxford (1993) [4].

## 3. COMPUTATIONAL PROCEDURES

The calculations were carried out using the GAUSSIAN 94 package of programs [5]. The computational method is Becke's gradient-corrected exchange-correlation density functionals (B3LYP) [6]. The unrestricted open-shell Hartree-Fock (UHF) formalism was applied to open-shell systems. According to the theorem of Hohenberg and Kohn [7], the functionals employed by DFT methods partition the electronic energy $E$ of a molecule into the terms

$$
E(\rho)=E^{\mathrm{T}}(\rho)+E^{\mathrm{V}}(\rho)+E^{\mathrm{J}}(\rho)+E^{\mathrm{XC}}(\rho)
$$

where $E^{\mathrm{T}}$ is the kinetic energy of the electrons, $E^{\mathrm{V}}$ is the potential energy of nuclear-electron attraction and nuclear-nuclear repulsion, $E^{\mathrm{J}}$ is the electron-electron repulsion of the classical energy of the density $\rho$, and $E^{\mathrm{XC}}$ is the exchange energy ( X ) arising from the wave function including the dynamical correlation (C) of electron motion. The term $E^{\mathrm{XC}}$ is divided into two separate functionals

$$
E^{\mathrm{XC}}(\rho)=E^{\mathrm{X}}(\rho)+E^{\mathrm{C}}(\rho)
$$

The definition of the functionals $E^{\mathrm{X}}(\rho)$ and $E^{\mathrm{C}}(\rho)$ can be found in the literature (Foresman, Frisch [8]). Becke introduced a gradient-corrected functional $E^{X}(\mathrm{~B})(\rho, \nabla \rho)$ and formulated functionals which include a mixture (hybrid) of Hartree-Fock (HF) exchange and DFT exchange (X) plus correlation (C) as

$$
E^{\mathrm{XC}}(\text { hybrid })=\mathrm{c}_{\mathrm{HF}} E^{\mathrm{X}}(\mathrm{HF})+\mathrm{c}_{\mathrm{DFT}} E^{\mathrm{XC}}(\mathrm{DFT})
$$

where the coefficients c are adjustable parameters. Becke's B3LYP functional, for instance, is a threeparameter functional of the following composition:

$$
\begin{aligned}
& E^{\mathrm{XC}}(\mathrm{~B} 3 \mathrm{LYP})=E^{\mathrm{X}}+\mathrm{c}_{0}\left[E^{\mathrm{X}}(\mathrm{HF})-E^{\mathrm{X}}(\mathrm{DFT})\right]+\mathrm{c}_{\mathrm{X}} E^{\mathrm{X}}(\mathrm{~B})+E^{\mathrm{C}}(\mathrm{VWN} 3)+ \\
& \mathrm{c}_{\mathrm{C}}\left[E^{\mathrm{C}}(\mathrm{LYP})-E^{\mathrm{C}}(\mathrm{VWN} 3)\right]
\end{aligned}
$$

where VWN is the Vosko, Wilk, Nusair functional [9], and LYP is the Lee, Yang, Parr functional [10]. The parameters $\mathrm{c}_{0}, \mathrm{c}_{\mathrm{X}}$, and $\mathrm{c}_{\mathrm{C}}$ are determined by fitting to atomization energies, ionization energies, proton affinities, and atomic energies of a set of molecules. Thus, the B3LYP procedure is semi-empirical in this sense. DFT calculations proceed in the same way as ab initio HF calculations, with the addition of the extra term $E^{\mathrm{XC}}$, which is computed via numerical integration.

A geometry optimization is complete when the force between the nuclei is below the cutoff value of $0.00045 E_{\mathrm{h}} a_{\mathrm{o}}^{-1}$, and the calculated displacement of the internuclear distance for the next optimization step is below $0.0018 a_{0}$. For weakly bound systems, scanning of the energy curve was applied in addition.

Vibrational wave numbers depend on second derivatives of the energy with respect to the nuclear positions. Analytic second derivatives are available for DFT calculations. The absolute absorption intensity $A$ which is measured in the unit of $\mathrm{km} \mathrm{mol}^{-1}$ is calculated by the formula

$$
A_{1-0}=\left(8 \pi^{3} / 3 \mathrm{hc}\right) N_{\mathrm{A}}\left|\mu_{1-0}\right|^{2} \omega_{1-0}
$$

where $N_{\mathrm{A}}$ is the Avogadro constant, $\mu_{1-0}$ is the electric dipole transition moment between the states 0 and 1 , and $\omega_{1-0}$ is the wave number. In the experimental literature, the quantity $S=A / R T$ which is measured in the unit of $\mathrm{cm}^{-2} \mathrm{~atm}^{-1}$ at a given temperature $T$ is used for gases. This quantity results from the substitution of the concentration $c$ in Beer's law by the partial pressure $p$. The conversion of units is given above.

## 4. BASIS SETS

The Gaussian atomic basis sets cc-pvqz (correlation consistent polarized valence quadruple-zeta) [11] are used for the atoms H to Ar. The description of the basis sets is arranged line by line in the following way: atom, basis set key word, standard notation of the basis set, contraction scheme of the basis set, size of the basis set.

| H | He | $\mathrm{Li}-\mathrm{Ne}$ |
| :--- | :--- | :--- |
| cc-pvqz | $\mathrm{cc}-\mathrm{pvqz}$ | $\mathrm{cc}-\mathrm{pvqz}$ |
| $4 \mathrm{~s} / 3 \mathrm{p} / 2 \mathrm{~d} / 1 \mathrm{f}$ | $4 \mathrm{~s} / 3 \mathrm{p} / 2 \mathrm{~d} / 1 \mathrm{f}$ | $5 \mathrm{~s} / 4 \mathrm{p} / 3 \mathrm{~d} / 2 \mathrm{f} / 1 \mathrm{~g}$ |
| $6.1 .1 .1 / 1.1 .1 / 1.1 / 1$ | $7.1 .1 .1 / 1.1 .1 / 1.1 / 1$ | $12.12 .1 .1 .1 / 6.1 .1 .1 / 1.1 .1 / 1.1 / 1$ |
| $35 / 30$ | $36 / 30$ | $92 / 55$ |

$\mathrm{Na}-\mathrm{Ar}$
cc-pvqz
$6 \mathrm{~s} / 5 \mathrm{p} / 3 \mathrm{~d} / 2 \mathrm{f} / 1 \mathrm{~g}$
16.16.16.1.1.1/11.11.1.1.1/1.1.1/1.1/1

164/59
Augmented basis sets (aug-cc-pvqz) were used for anionic systems of the type $\mathrm{A}^{-}$(atomic anions), the results of which are listed in Table 1, and $\mathrm{AH}^{-}$(diatomic hydrides). The key word aug-cc-pvqz results in additional diffuse functions of the type one s, one p set, one d set, etc. Augmented basis sets are not mentioned in Table 3 for diatomic molecules; instead, their use for $E_{\mathrm{i}}\left(\mathrm{AH}^{-}\right)$is emphasized here. Augmented basis sets are also necessary, according to Table 1 for atoms, for $D_{0}\left(\mathrm{AB}^{-}\right)$where one of the dissociating atoms is either $\mathrm{O}^{-}$or $\mathrm{F}^{-}$.

Table 1 Total electronic energies $E / E_{\mathrm{h}}$ and ionization energies $E_{\mathrm{i}} / \mathrm{eV}$ of the atoms $\mathrm{H}-\mathrm{Ar}$ and their stable anions.

| Atom | State | E(cc-pvqz) | $E$ (aug-cc-pvqz) | $E_{\mathrm{i}}($ calc. $)$ | $E_{\mathrm{i}}(\mathrm{exp}$. |
| :---: | :---: | :---: | :---: | :---: | :---: |
| H | ${ }^{2} \mathrm{~S}$ | -0.502346* | -0.502391 | 13.67 | 13.61 |
| $\mathrm{H}^{-}$ | ${ }^{1} \mathrm{~S}$ | -0.517503 | -0.535949* | 0.91 | 0.75 |
| He | ${ }^{1} \mathrm{~S}$ | -2.914981 |  | 24.94 | 24.59 |
| $\mathrm{He}^{+}$ | ${ }^{2} \mathrm{~S}$ | -1.998484 |  |  |  |
| Li | ${ }^{2} \mathrm{~S}$ | -7.492222 |  | 5.62 | 5.39 |
| $\mathrm{Li}^{+}$ | ${ }^{1} \mathrm{~S}$ | -7.285508 |  |  |  |
| $\mathrm{Li}^{-}$ | ${ }^{1} \mathrm{~S}$ | -7.509870 |  | 0.48 | 0.62 |
| Be | ${ }^{1} \mathrm{~S}$ | -14.672628 |  | 9.11 | 9.32 |
| $\mathrm{Be}^{+}$ | ${ }^{2} \mathrm{~S}$ | -14.337861 |  |  |  |
| $\mathrm{Be}^{-}$ | ${ }^{2} \mathrm{P}$ | -14.654534 |  | - | - |
| B | ${ }^{2} \mathrm{P}$ | -24.664786 |  | 8.74 | 8.30 |
| $\mathrm{B}^{+}$ | ${ }^{1} \mathrm{~S}$ | -24.343692 |  |  |  |
| $\mathrm{B}^{-}$ | ${ }^{3} \mathrm{P}$ | -24.670391 |  | 0.15 | 0.28 |
| C | ${ }^{3} \mathrm{P}$ | -37.860591* | -37.860785 | 11.54 | 11.26 |
| $\mathrm{C}^{+}$ | ${ }^{2} \mathrm{P}$ | -37.436638 |  |  |  |
| $\mathrm{C}^{-}$ | ${ }^{4} \mathrm{~S}$ | -37.898495 | -37.911284* | 1.37 | 1.27 |
| N | ${ }^{4} \mathrm{~S}$ | -54.605328* | -54.605735 | 14.65 | 14.54 |
| $\mathrm{N}^{+}$ | ${ }^{3} \mathrm{P}$ | -54.066941 |  |  |  |
| O | ${ }^{3} \mathrm{P}$ | -75.098193* | -75.099041 | 14.12 | 13.61 |
| $\mathrm{O}^{+}$ | ${ }^{4} \mathrm{~S}$ | -74.579257 |  |  |  |
| $\mathrm{O}^{-}$ | ${ }^{2} \mathrm{P}$ | -75.136450 | -75.160853* | 1.68 | 1.46 |
| F | ${ }^{2} \mathrm{P}$ | -99.772525* | -99.773643 | 17.70 | 17.42 |
| (Continued on next page) |  |  |  |  |  |

Table 1 (Continued)

| Atom | State | $E$ (cc-pvqz) | $E$ (aug-cc-pvqz) | $E_{\mathrm{i}}($ calc. $)$ | $E_{\mathrm{i}}($ exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{F}^{+}$ | ${ }^{3} \mathrm{P}$ | -99.121972 |  |  |  |
| $\mathrm{F}^{-}$ | ${ }^{1} \mathrm{~S}$ | -99.877807 | -99.903238* | 3.53 | 3.40 |
| Ne | ${ }^{1} \mathrm{~S}$ | -128.975664* | -128.977143 | 21.67 | 21.56 |
| $\mathrm{Ne}^{+}$ | ${ }^{2} \mathrm{P}$ | -128.179076 |  |  |  |
| Na | ${ }^{2} \mathrm{~S}$ | -162.298962 |  | 5.43 | 5.14 |
| $\mathrm{Na}^{+}$ | ${ }^{1} \mathrm{~S}$ | -162.099313 |  |  |  |
| $\mathrm{Na}^{-}$ | ${ }^{1} \mathrm{~S}$ | -162.319312 |  | 0.58 | 0.55 |
| Mg | ${ }^{1} \mathrm{~S}$ | -200.098962 |  | 7.73 | 7.65 |
| $\mathrm{Mg}^{+}$ | ${ }^{2} \mathrm{~S}$ | -199.815032 |  |  |  |
| $\mathrm{Mg}^{-}$ | ${ }^{2} \mathrm{P},{ }^{2} \mathrm{~S}$ | -200.085222 |  | - | - |
| Al | ${ }^{2} \mathrm{P}$ | -242.393290* | -242.393323 | 6.02 | 5.99 |
| $\mathrm{Al}^{+}$ | ${ }^{1} \mathrm{~S}$ | -242.171973 |  |  |  |
| $\mathrm{Al}^{-}$ | ${ }^{3} \mathrm{P}$ | -242.403407 | -242.410353* | 0.46 | 0.44 |
| Si | ${ }^{3} \mathrm{P}$ | -289.399119* | -289.399169 | 8.11 | 8.15 |
| $\mathrm{Si}^{+}$ | ${ }^{2} \mathrm{P}$ | -289.100932 |  |  |  |
| $\mathrm{Si}^{-}$ | ${ }^{4} \mathrm{~S}$ | -289.443146 | -289.448567* | 1.34 | 1.39 |
| P | ${ }^{4} \mathrm{~S}$ | -341.288711* | -341.288791 | 10.39 | 10.49 |
| $\mathrm{P}^{+}$ | ${ }^{3} \mathrm{P}$ | -340.907043 |  |  |  |
| $\mathrm{P}^{-}$ | ${ }^{3} \mathrm{P}$ | -341.315305 | -341.324065* | 0.96 | 0.75 |
| S | ${ }^{3} \mathrm{P}$ | -398.142105* | -398.142260 | 10.55 | 10.36 |
| $\mathrm{S}^{+}$ | ${ }^{4} \mathrm{~S}$ | -397.754305 |  |  |  |
| $\mathrm{S}^{-}$ | ${ }^{2} \mathrm{P}$ | -398.215963 | -398.223216* | 2.20 | 2.08 |
| Cl | ${ }^{2} \mathrm{P}$ | -460.178519* | -460.178678 | 13.06 | 12.97 |
| $\mathrm{Cl}^{+}$ | ${ }^{3} \mathrm{P}$ | -459.698367 |  |  |  |
| $\mathrm{Cl}^{-}$ | ${ }^{1} \mathrm{~S}$ | -460.307494 | -460.313669* | 3.67 | 3.62 |
| Ar | ${ }^{1} \mathrm{~S}$ | -527.563508* | -527.563658 | 15.79 | 15.76 |
| $\mathrm{Ar}^{+}$ | ${ }^{2} \mathrm{P}$ | -526.983284 |  |  |  |

*Total energy used for $E_{\mathrm{i}}$

## 5. STATISTICS OF DEVIATIONS BETWEEN CALCULATED AND EXPERIMENTAL PROPERTIES OF DIATOMIC MOLECULES

For $N$ comparisons of calculated with experimental values of a property $p$, the mean absolute deviation

$$
\overline{|\Delta p|}=\Sigma_{i}^{N}\left|\Delta p_{i}\right| / N=\Sigma_{i}^{N} \mid p_{i}(\text { calc. })-p_{i}(\text { (exp. }) \mid / N
$$

the average value

$$
\bar{p}=\Sigma_{i}^{N} p_{i} / N
$$

and the mean relative deviation
$\overline{|\Delta p|} / \bar{p}$
are listed in Table 2.
The statistics for equilibrium internuclear distances $r_{\mathrm{e}}$ and harmonic vibrational wave numbers $\omega_{\mathrm{e}}$ incorporates data for electronic ground states. The statistics for $r_{\mathrm{e}}$ and $\omega_{\mathrm{e}}$, presented in Table 2, are more successful than those for ab initio HF and MP2 (Moeller-Plesset second-order perturbation theory) calculations of the literature (Hehre et al. [12]). Statistics of the past, however, suffer from limited numbers of comparisons and also from the omission of worst cases. Even more successful are the present

Table 2 Statistics for properties of diatomic molecules. $N$ : number of comparisons between calculated and experimental values of property $p ; \overline{|\Delta p|}$ : mean absolute deviation; $\bar{p}$ : average value of calculated data; $\overline{\Delta p \mid} / \bar{p}$ : mean relative deviation; $|\Delta p|$ (max.): maximum absolute deviation between calculated and experimental value of property $p$.

| $\begin{aligned} & \hline N \\ & 138 \end{aligned}$ | $\begin{aligned} & \hline \Delta r_{e} \mid \\ & 0.0125 \AA \end{aligned}$ | $\begin{aligned} & \overline{r_{e}} \\ & 1.5994 \AA \end{aligned}$ | $\begin{aligned} & \overline{\Delta r_{e}} / / \overline{r_{e}} \\ & 0.0078 \end{aligned}$ | $\begin{aligned} & \left\|\Delta r_{e}\right\|(\max .) \\ & 0.0645 \AA\left(\mathrm{He}_{2}{ }^{+}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & N \\ & 133 \end{aligned}$ | $\begin{aligned} & \mid \overline{\Delta \omega_{e} \mid} \\ & 40 \mathrm{~cm}^{-1} \end{aligned}$ | $\begin{aligned} & \overline{\omega_{e}} \\ & 1430 \mathrm{~cm}^{-1} \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta \omega_{e}}\right\| / \overline{\omega_{e}} \\ & 0.0282 \end{aligned}$ | $\begin{aligned} & \left\|\Delta \omega_{e}\right\|(\max .) \\ & 338 \mathrm{~cm}^{-1}\left(\mathrm{He}_{2}{ }^{+}\right) \end{aligned}$ |
| $N$ 40 | $\begin{aligned} & \overline{\Delta \mu \mid} \\ & 0.17 \mathrm{D} \end{aligned}$ | $\begin{aligned} & \bar{\mu} \\ & 2.35 \mathrm{D} \end{aligned}$ | $\begin{aligned} & \|\overline{\Delta \mu \mid}\| / \bar{\mu} \\ & 0.0746 \end{aligned}$ | $\begin{aligned} & \|\Delta \mu\|(\max .) \\ & 0.94 \mathrm{D}(\mathrm{NaH}) \end{aligned}$ |
| $\begin{aligned} & N \\ & 50^{\mathrm{a}} \\ & 5^{\mathrm{b}} \\ & 38^{\mathrm{c}} \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta E_{i}}\right\| \\ & 0.19 \mathrm{eV} \\ & 0.78 \mathrm{eV} \\ & 0.17 \mathrm{eV} \end{aligned}$ | $\begin{aligned} & \overline{E_{i}} \\ & 10.41 \mathrm{eV} \\ & 29.46 \mathrm{eV} \\ & 1.41 \mathrm{eV} \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta E_{i}}\right\| / \overline{E_{i}} \\ & 0.0179 \\ & 0.0265 \\ & 0.1204 \end{aligned}$ | $\begin{aligned} & \left\|\Delta E_{i}\right\|(\max .) \\ & 1.02 \mathrm{eV}(\mathrm{LiMg}) \\ & 2.08 \mathrm{eV}\left(\mathrm{~N}_{2}^{+}\right) \\ & 0.70 \mathrm{eV}\left(\mathrm{C}_{2}^{-}\right) \end{aligned}$ |
| $\begin{aligned} & N \\ & 155 \end{aligned}$ | $\begin{aligned} & \mid \overline{\left\|\Delta D_{0}\right\|} \\ & 0.17 \mathrm{eV} \end{aligned}$ | $\begin{aligned} & \overline{D_{0}} \\ & 3.86 \mathrm{eV} \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta D_{0}}\right\| / \overline{D_{0}} \\ & 0.0433 \end{aligned}$ | $\begin{aligned} & \left\|\Delta D_{0}\right\|(\text { max. }) \\ & 1.23 \mathrm{eV}\left(\mathrm{Ne}_{2}^{+}\right) \end{aligned}$ |
| $\begin{aligned} & N \\ & 31 \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta \omega_{e} x_{e}}\right\| \\ & 2.56 \mathrm{~cm}^{-1} \end{aligned}$ | $\begin{aligned} & \omega_{e} x_{e} \\ & 37.8 \mathrm{~cm}^{-1} \end{aligned}$ | $\begin{aligned} & \left\|\overline{\Delta \omega_{e} x_{e}}\right\| / \overline{\omega_{e} x_{e}} \\ & 0.0677 \end{aligned}$ | $\begin{aligned} & \left\|\Delta \omega_{e} x_{e}\right\|(\text { max. }) \\ & 7.35 \mathrm{~cm}^{-1}(\mathrm{NH}) \end{aligned}$ |

${ }^{a}$ neutral molecules, ${ }^{\mathrm{b}}$ cations, ${ }^{\mathrm{c}}{ }^{\text {anions }}$
statistics for dissociation energies $D_{0}$ which overrule in some sense the skepticism as exemplified by the remark of Benson [13] in 1981: "Ab initio methods for solving the Schrödinger equation are not yet capable of providing results of $\Delta_{\mathrm{f}} H$ to better than $\pm 4 \mathrm{kcal} \mathrm{mol}^{-1}$ even for relatively simple species and the prospect is not bright for any major breakthroughs in this area in the next decade or so". (See for comparison in Table 2: $\left.\overline{\Delta D_{0} \mid}=0.17 \mathrm{eV}=3.8 \mathrm{kcal} \mathrm{mol}^{-1}\right)^{*}$. No statistics are presented for absorption intensities $A$ owing to the limited number of only five comparisons.

In view of the combined theoretical and experimental uncertainties, the following aspect should be taken into account. If there is a discrepancy between a calculated and experimental value, there is always a possibility that it is the latter that is in error. Famous resolved examples are $\mathrm{CH}^{-}$and $\mathrm{SiH}^{-}$, or SiCl and SiF . The cation $\mathrm{HeNe}^{+}$, for instance, which is excluded from the statistics, is an unresolved example where data from calculations [B3LYP and $\operatorname{CCSD}(\mathrm{T})$ ] differ dramatically from uncertain experimental data. Uncertain experimental properties are set in parentheses and are generally not considered in the present statistics.

Statistics can be found in the literature where the G2 procedure was employed. G2 is not a computational method but a series of ab initio calculations, including empirical corrections, to use the additivity of basis set extension and the extension of electron correlation for the total electronic energy. The application is, therefore, restricted to the properties $E_{\mathrm{i}}$ and $D_{0}$. The limited number of diatomic systems, $25\left(E_{\mathrm{i}}\right)$ and $26\left(D_{0}\right)$, makes statistics questionable [14]. The G2 procedure was applied also to small molecules, $42\left(E_{\mathrm{i}}\right)$ and $56\left(D_{0}\right)$. For both properties, mean absolute deviations of 0.05 eV were reported. However, worst cases were omitted, such as $\mathrm{C}_{2}{ }^{-}(0.44), \mathrm{CN}(0.42), \operatorname{LiMg}(1.23)$ for $E_{\mathrm{i}}$ (absolute deviations in eV in parentheses), and $\mathrm{CCl}(0.82), \mathrm{MgO}(0.96)$ for $D_{0}$ (absolute deviations in eV in parentheses) [15]. A variety of $a b$ initio methods was applied to 120 first- and second-row diatomic molecules where the properties $T_{\mathrm{e}}, r_{\mathrm{e}}, \omega_{\mathrm{e}}$, and $D_{0}$ are presented; however, ions as well as rare gas compounds were omitted [16].

[^0]Recently performed statistics from the literature for $r_{\mathrm{e}}, \omega_{\mathrm{e}}, E_{\mathrm{i}}$, and $D_{0}$ still suffer from the limited sets of data, but the comparison with different DFT methods globally favors the B3LYP method [17]. The G2 procedure for the calculation of $E_{\mathrm{i}}$ and $D_{0}$ was very recently replaced by G3 where the B3LYP method is involved [18]. No statistics could be found in the literature for anharmonicities, $\omega_{\mathrm{e}} x_{\mathrm{e}}$, NMR chemical shifts, $\delta$, for the most abundant nuclei, dipole moments, $\mu$, and isotropic hyperfine coupling constants $A_{\text {iso }}$.

## 6. TABLE OF PROPERTIES OF DIATOMIC MOLECULES

Table 3 for properties of diatomic molecules is arranged as follows. Calculated data can be found in the lines with the boldface molecular formulae on the far left. These lines have a blank space in the far right column, indicated by Ref.(Exp.). The corresponding following lines with one or more reference numbers on the outer right show experimental values. In some cases, electronically excited states were calculated, in addition to the ground state. One line is used per state for calculated as well as experimental data (cf., for example, CH ). Sometimes, in particular for cations $\mathrm{AB}^{+}$, there is a full blank line instead of experimental values. In such a case, the ionization energy $\left(E_{\mathrm{i}}\right)$ of the corresponding neutral compound AB is experimentally known, and therefore, the cation $\mathrm{AB}^{+}$is regarded as an identified diatomic system (cf., for example, $\mathrm{AlF}^{+}$), irrespective of additional known properties. Uncertain experimental properties are set in parentheses (cf., for example, $\mathrm{Cl}_{2}^{-}$).

Table 3 B3LYP calculated and experimental properties of diatomic molecules for $\mathrm{H}-\mathrm{Ar} . T_{\mathrm{e}} / \mathrm{cm}^{-1}$ : relative electronic energy referred to the ground state; $r_{\mathrm{e}} / \AA$ : equilibrium internuclear distance; $\omega_{\mathrm{e}} / \mathrm{cm}^{-1}$ : harmonic vibrational wavenumber; $A / \mathrm{km} \mathrm{mol}^{-1}$ : vibrational absolute absorption intensity; $\mu / \mathrm{D}$ : electric dipole moment; $E_{\mathrm{i}} / \mathrm{eV}$ : adiabatic, zero-point energy corrected ionization energy; $D_{0} / \mathrm{eV}$ : dissociation energy, corrected for zeropoint energy. For conversion of units see chapter 2.

| Molecule | State | $T_{\mathrm{e}}$ | $r_{\text {e }}$ | $\omega_{\mathrm{e}}$ | A | $\mu$ | $E_{\mathrm{i}}$ | $D_{0}$ | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{27} \mathrm{Al}_{2}$ | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{-}$ | 0 | 2.5020 | 324 | 0 | 0 | 1.20 |  | 1,19 |
|  | ${ }^{3} \Sigma^{8}$ - | 0 | 2.466 | 350 |  |  |  | 1.34 |  |
| ${ }^{27} \mathrm{Al}_{2}{ }^{-}$ | ${ }^{2} \Sigma_{\text {u }}{ }^{\text {- }}$ | 0 | 2.4493 | 345 | 0 | 0 | 0.81 | 1.74 |  |
|  |  |  |  |  |  |  | 1.10 |  | 19 |
| ${ }^{27} \mathrm{Al}^{35} \mathrm{Cl}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.1619 | 458 | 126 | -1.77 | 9.26 | 5.04 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 2.1301 | 481 |  | $1.5 \pm 0.5$ | 9.4 | 5.26 | 1,19 |
| ${ }^{27} \mathrm{Al}^{35} \mathrm{Cl}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 2.0368 | 577 | 60 | -5.62 | $1.81\left(\mathrm{Al}^{+}+\mathrm{Cl}\right)$ |  |  |
|  |  |  |  |  |  |  |  |  | 19 |
| ${ }^{27} \mathrm{Al}^{19} \mathrm{~F}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.6735 | 780 | 108 | -1.55 | 9.73 | 6.88 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.6544 | 802 |  | 1.53 | 9.73 | 6.84 | 1,19 |
| ${ }^{27} \mathrm{Al}^{19} \mathrm{~F}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.6113 | 925 | 79 | -5.35 |  | $3.18\left(\mathrm{Al}^{+}+\mathrm{F}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 1 |
| ${ }^{27} \mathrm{Al}^{1} \mathrm{H}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.6599 | 1653 | 704 | -0.20 | 8.23 | 3.06 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.6474 | 1682 |  | 0.32 |  | 2.91 | 19,20 |
| ${ }^{27} \mathbf{A l ~}^{1} \mathbf{H}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.6198 | 1632 | 22 | -0.62 |  | $0.85\left(\mathrm{Al}^{+}+\mathrm{H}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.6018 | 1620 |  |  |  |  | 1 |
| ${ }^{27} \mathrm{Al}^{14} \mathrm{~N}$ | ${ }^{3} \Pi$ | 0 | 1.7920 | 748 | 9 | -2.80 |  | 2.73 | 1 |
|  | ${ }^{3} \Pi$ | 0 | 1.7864 | 747 |  |  |  | $3.04 \pm 1.0$ | 1,19 |
| ${ }^{27} \mathrm{Al}^{16} \mathrm{O}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.6290 | 955 | 1 | -4.21 | 9.86 | 5.16 |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.6179 | 979 |  |  | 9.53 | 5.27 | 1 |
| ${ }^{27} \mathrm{Al}^{16} \mathrm{O}^{+}$ | ${ }^{1} \Sigma^{+}$ | 3815 | 1.6035 | 1008 | 2 | -7.17 |  |  |  |
|  | ${ }^{3} \Pi$ | 0 | 1.7459 | 738 | 12 | -4.59 |  | $1.32\left(\mathrm{Al}^{+}+\mathrm{O}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 1 |
| ${ }^{27} \mathrm{Al}^{16} \mathrm{O}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.6425 | 981 | 59 | -1.16 | $\begin{aligned} & 2.44 \\ & 2.60 \end{aligned}$ | $6.55\left(\mathrm{Al}+\mathrm{O}^{-}\right)$ | 19 |
| ${ }^{27} \mathrm{Al}^{31} \mathbf{P}$ | ${ }^{3} \Pi$ | 0 | 2.2234 | 456 | 8 | -2.48 |  | 2.25 |  |

Table 3 (Continued)


Table 3 (Continued)

(Continued on next page)

Table 3 (Continued)

(Continued on next page)

Table 3 (Continued)

| Molecule | State | $T_{\mathrm{e}}$ | $r_{\text {e }}$ | $\omega_{\text {e }}$ | A | $\mu$ | $E_{\mathrm{i}}$ | $D_{0}$ | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ${ }^{1} \Sigma^{+} \quad 0$ |  |  |  |  |  |  |  | 37 |
| ${ }^{4} \mathrm{He}^{40} \mathrm{Ar}$ |  |  | (repulsive potential) |  |  |  |  |  |  |
|  | $\begin{aligned} & { }^{1} \Sigma^{+} \\ & { }^{2} \Sigma^{+} \end{aligned}$ | 0 | 3.51 |  | 24 | +0.19 |  | $0.0024\left(D_{\mathrm{e}}\right)$ | 1 |
| ${ }^{4} \mathrm{He}^{40} \mathrm{Ar}^{+}$ |  | 0 | 2.1236 | 329 |  |  |  | $0.12\left(\mathrm{He}+\mathrm{Ar}^{+}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 1 |
| ${ }^{4} \mathrm{He}^{\mathbf{1}} \mathrm{H}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 0.7879 | 3073 | 776 | $\begin{aligned} & +1.27 \\ & (1.66) \end{aligned}$ |  | $1.82\left(\mathrm{He}+\mathrm{H}^{+}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 0.7743 | 2911 |  |  |  | 1.84 | 1,23 |
| ${ }^{4} \mathrm{He}^{20} \mathrm{Ne}$ | ${ }^{1} \Sigma^{+}$ | 0 | (repulsive potential) |  |  |  |  |  |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 3.21 |  |  |  |  |  | 0.0002 | 1 |
| ${ }^{4} \mathrm{He}^{\mathbf{2 0}} \mathrm{Ne}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.4970 | 808 | 0 | -1.37 |  | $1.50\left(\mathrm{He}+\mathrm{Ne}^{+}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | (1.300) | (1308) |  |  |  | (0.69) | 1 |
| ${ }^{1} \mathbf{H}^{19} \mathrm{~F}$ | ${ }^{1} \Sigma^{+}$ | 0 | 0.9214 | 4088 | 104 | -1.82 | 16.01 | 5.76 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 0.9168 | 4138 |  | 1.83 | 16.04 | 5.87 | 1 |
| ${ }^{1} \mathbf{H}^{19} \mathbf{F}^{+}$ | ${ }^{2} \Pi$ | 0 | 1.0133 | 2986 | 541 | -2.36 |  | $3.42\left(\mathrm{H}^{+}+\mathrm{F}\right)$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.0011 | 3090 |  |  |  | 3.42 | 1 |
| ${ }^{7} \mathrm{Li}_{2}$ | $\begin{aligned} & { }^{1} \Sigma^{+}{ }^{+}{ }^{g_{+}}{ }^{+} \end{aligned}$ | 0 | 2.6910 | 345 | 0 | 0 | 5.23 | 0.98 |  |
|  |  | 0 | 2.6729 | 351 |  |  | 5.14 | 1.05 | 1,38 |
| ${ }^{7} \mathrm{Li}_{2}{ }^{+}$ | ${ }^{2} \Sigma_{\mathrm{g}}{ }^{\text {a }}$ | 0 | 3.0623 | 251 | 0 | 0 |  | 1.28 |  |
|  |  |  | 3.12 | 262 |  |  |  | 1.30 | 38 |
| ${ }^{7} \mathrm{Li}_{2}{ }^{-}$ | ${ }^{2} \Sigma_{u}+$ | 0 | 3.0311 | 225 | 0 | 0 | 0.43 | 0.83 |  |
|  |  |  |  |  |  |  |  | (0.88) | 1 |
| ${ }^{7} \mathbf{L i}{ }^{\mathbf{2 7}} \mathbf{A l}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.8794 | 306 | 12 | -2.66 |  | 0.89 |  |
|  |  |  |  |  |  |  |  | 0.75 | 19 |
| ${ }^{7} \mathbf{L i}^{40} \mathrm{Ar}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.3446 | 283 | 150 | -8.25 |  | $0.31\left(\mathrm{Li}^{+}+\mathrm{Ar}\right)$ |  |
|  |  |  |  |  |  |  |  | $0.30\left(D_{\mathrm{e}}\right)$ | 1 |
| ${ }^{7} \mathbf{L i}{ }^{9} \mathrm{Be}$ | ${ }^{2} \Pi$ | 4919 | 2.2328 | 502 | 40 | -3.73 |  |  |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 2.5600 | 351 | 11 | -2.81 |  | 0.45 |  |
|  | ${ }^{2} \Sigma^{+}$ |  | $2.59\left(r_{0}\right)$ | 295( $\omega_{0}$ ) |  |  |  |  | 39 |
| ${ }^{7} \mathrm{Li}^{35} \mathrm{Cl}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.0200 | 643 | 121 | -6.90 | 9.80 | 4.73 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 2.0207 | 643 |  | 7.13 | 9.57 | $4.82 \pm 0.1$ | 1,19 |
| ${ }^{7} \mathbf{L i}^{\mathbf{3 5}} \mathrm{Cl}^{+}$ | ${ }^{2} \Sigma^{+}$ | 4716 | 2.7873 | 109 | 163 | -10.08 |  |  |  |
|  | ${ }^{2} \Pi$ | 0 | 2.3823 | 325 | 135 | -7.97 | $0.55\left(\mathrm{Li}^{+}+\mathrm{Cl}\right)$ |  |  |
|  |  |  |  |  |  |  |  |  | 19 |
| ${ }^{7} \mathbf{L i}^{35} \mathrm{Cl}^{-}$ | ${ }^{2} \Sigma^{+}$ | 0 | 2.1300 | 501 | 25 | +7.35 | 0.63 | $1.85\left(\mathrm{Li}+\mathrm{Cl}^{-}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 2.123 | 500 |  |  | 0.59 | 1.75 | 40 |
| ${ }^{7} \mathbf{L i}^{19} \mathbf{F}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.5609 | 915 | 140 | -6.12 |  | 5.90 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.5639 | 910 |  | 6.28 |  | $5.94 \pm 0.2$ | 1,19 |
| ${ }^{7} \mathbf{L i}^{1} \mathrm{H}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.5900 | 1406 | 173 | -5.70 | 7.93 | 2.45 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.5957 | 1406 |  | 5.88 | 7.7 | 2.43 | 1,19 |
| ${ }^{7} \mathbf{L i}^{1} \mathbf{H}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 2.1801 | 468 | 20 | -1.91 |  | $0.16\left(\mathrm{Li}^{+}+\mathrm{H}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 19 |
| ${ }^{7} \mathbf{L i}^{\mathbf{1}} \mathbf{H}^{-}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.6865 | 1145 | 715 | +3.36 | 0.35 | $2.40\left(\mathrm{Li}+\mathrm{H}^{-}\right)$ |  |
|  |  |  |  |  |  |  | 0.34 |  | 19 |
| ${ }^{7} \mathbf{L i}^{4} \mathrm{He}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.8721 | 309 | 60 | -3.31 |  | $0.08\left(\mathrm{Li}^{+}+\mathrm{He}\right)$ |  |
|  |  |  |  |  |  |  |  | 0.07 | 32 |
| ${ }^{7} \mathrm{Li}^{\mathbf{2 4}} \mathrm{Mg}$ | ${ }^{2} \Sigma^{+}$ | 0 | 3.0954 | 186 | 1 | -0.57 | 4.94 | 0.18 |  |
|  | ${ }^{2} \Sigma^{+}$ |  |  |  |  |  | 5.96 | 0.66 | 19,41 |
| ${ }^{7} \mathbf{L i}^{\mathbf{2 4}} \mathbf{M g}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.8846 | 267 | 93 | -5.68 |  | $0.86\left(\mathrm{Li}^{+}+\mathrm{Mg}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 41 |
| ${ }^{7} \mathrm{Li}^{14} \mathrm{~N}$ | ${ }^{3} \Pi$ | 2015 | 1.7114 | 792 | 123 | -5.70 |  |  |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.8542 | 679 | 18 | -6.64 |  | 1.89 |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 |  |  |  |  |  |  | 2 |
| ${ }^{7} \mathrm{Li}^{23} \mathrm{Na}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.8758 | 254 | 0 | $+0.02$ | 5.17 | 0.80 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 2.815 | 257 |  | 0.46 | 5.05 | 0.86 | 1,19 |
| ${ }^{7} \mathrm{Li}^{\mathbf{2 3}} \mathrm{Na}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 3.3361 | 180 | 9 | $-4.32$ |  | $1.05\left(\mathrm{Li}+\mathrm{Na}^{+}\right)$ |  |
|  |  |  |  |  |  |  |  | (Continued | next page) |

Table 3 (Continued)

| Molecule | State | $T_{\mathrm{e}}$ | $r_{\text {e }}$ | $\omega_{\text {e }}$ | A | $\mu$ | $E_{\mathrm{i}}$ | $D_{0}$ | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{7} \mathrm{Li}^{16} \mathrm{O}$ |  |  |  |  |  |  |  | 0.99 | 1 |
|  | ${ }^{2} \Sigma^{+}$ | 2290 | 1.5858 | 871 | 166 | -5.65 |  |  |  |
|  | ${ }^{2} \Pi$ | 0 | 1.6821 | 821 | 99 | -6.39 | 8.80 | 3.70 |  |
|  | ${ }^{2} \Sigma^{+}$ | 2565 |  |  |  |  |  |  | 42 |
|  | ${ }^{2} \Pi$ | 0 | 1.6882 | 815 |  | 6.84 | 8.45 | $3.42 \pm 0.1$ | 1,19,42 |
| ${ }^{7} \mathbf{L i}^{16} \mathrm{O}^{+}$ | ${ }^{3} \Sigma^{-}$ | 0 | 2.0736 | 365 | 94 | -6.41 |  | $0.53\left(\mathrm{Li}^{+}+\mathrm{O}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 1 |
| ${ }^{7} \mathbf{L i}{ }^{31} \mathbf{P}$ | ${ }^{3} \Sigma^{-}$ | 0 | 2.3249 | 490 | 25 | -6.57 |  | 1.86 |  |
|  | ${ }^{3} \Sigma^{-}$ |  |  |  |  |  |  | $2.53 \pm 0.24$ | 41 |
| ${ }^{7} \mathbf{L i}{ }^{32} \mathbf{S}$ | ${ }^{2} \Pi$ | 0 | 2.1500 | 576 | 81 | -6.86 |  | 3.17 |  |
|  |  |  |  |  |  |  |  | 3.20; $3.37 \pm 0.4$ | 41 |
| ${ }^{7} \mathbf{L i}{ }^{\mathbf{2 8}} \mathbf{S i}$ | ${ }^{4} \Sigma^{-}$ | 0 | 2.3629 | 448 | 57 | -5.79 |  | 1.63 |  |
|  |  |  |  |  |  |  |  | 1.54 | 41 |
| ${ }^{24} \mathrm{Mg}_{2}$ | ${ }^{1} \Sigma_{\mathrm{g}}^{+}$ | 0 |  | 47 | 0 | 0 |  | 0.01 |  |
|  | ${ }^{1} \Sigma_{\mathrm{g}}^{\mathrm{g}_{+}}$ | 0 | 3.8905 | 51 |  |  |  | 0.05 | 1 |
| ${ }^{24} \mathrm{Mg}^{35} \mathrm{Cl}$ | ${ }^{2} \Sigma^{\text {g }}$ | 0 | 2.2211 | 447 | 70 | -3.54 |  | 3.14 |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 2.1964 | 462 |  |  |  | 3.36 | 1,19,43 |
| ${ }^{24} \mathrm{Mg}^{35} \mathrm{Cl}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.4091 | 286 | 69 | +3.54 | 1.57 | $1.20\left(\mathrm{Mg}+\mathrm{Cl}^{-}\right)$ | 1,19,4 |
|  |  |  |  |  |  |  | 1.59 |  | 19 |
| ${ }^{24} \mathrm{Mg}^{19} \mathrm{~F}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.7659 | 706 | 72 | -3.26 |  | 4.45 |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.7500 | 720 |  |  |  | 4.75 | 1 |
| ${ }^{24} \mathrm{Mg}^{\mathbf{1}} \mathrm{H}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.7400 | 1456 | 245 | -1.37 | 7.11 | 1.36 |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.7297 | 1495 |  |  |  | 1.34 | 1 |
| ${ }^{24} \mathrm{Mg}^{1} \mathrm{H}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.6546 | 1722 | 17 | -3.51 |  | $1.97\left(\mathrm{Mg}^{+}+\mathrm{H}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.6519 | 1699 |  |  |  | 2.08 | 1 |
| ${ }^{24} \mathbf{M g}^{1} \mathrm{H}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.8899 | 1070 | 1928 | +1.74 | 0.81 | $1.26\left(\mathrm{Mg}+\mathrm{H}^{-}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 |  |  |  |  | 1.05 | 1.53 | 27 |
| ${ }^{24} \mathbf{M g}^{16} \mathrm{O}$ | ${ }^{1} \Sigma^{+}$ | 1089 | $1.7389$ |  | 6 | $-7.19$ |  |  |  |
|  | ${ }^{3} \Pi$ | 0 | 1.8882 | 626 | 42 | $-3.11$ |  | 2.37 |  |
|  | ${ }^{3} \Pi$ | 2623 | 1.870 | 648 |  |  |  |  | 44 |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.7490 | 785 |  |  |  | $2.56 \pm 0.2$ | 1,44 |
| ${ }^{24} \mathrm{Mg}^{32} \mathrm{~S}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.1492 | 528 | 6 | -7.46 |  | 1.99 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 2.1425 | 529 |  |  |  | 2.39 | 1,19 |
| ${ }^{14} \mathrm{~N}_{2}$ | ${ }^{1} \Sigma_{\mathrm{g}}{ }^{+}$ | 0 | 1.0902 | 2447 | 0 | 0 | 15.80 | 9.82 |  |
|  | ${ }^{1} \Sigma_{\mathrm{g}}{ }^{\text {+ }}$ | 0 | 1.0977 | 2358 |  |  | 15.58 | 9.76 | 1 |
| ${ }^{14} \mathrm{~N}_{2}{ }^{+}$ | ${ }^{2} \Sigma_{\mathrm{g}}{ }^{\text {g }}$ | 0 | 1.1035 | 2330 | 0 | 0 | 29.18 | 8.66 |  |
|  |  | 0 | 1.1164 | 2207 |  |  | 27.1 | 8.71 | 1 |
| ${ }^{14} \mathrm{~N}_{2}{ }^{++}$ | ${ }^{1} \Sigma_{\mathrm{g}}^{\mathrm{g}}{ }^{+}$ | 12281 | 1.1327 | 2077 | 0 | 0 |  |  |  |
|  | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{\text {- }}$ | 2288 | 1.3517 | 1192 | 0 | 0 |  |  |  |
|  | ${ }^{3} \square_{\text {I }}^{\text {u }}$ | 0 | 1.2186 | 1654 | 0 | 0 |  | $-4.32\left(\mathrm{~N}^{+}+\mathrm{N}^{+}\right)$ |  |
|  |  | 880 |  | 1460 |  |  |  | $-4.32\left({ }^{+}+\mathrm{N}\right)$ | 45 |
|  | ${ }^{1} \Sigma_{\mathrm{g}}{ }^{+}$ | 0 | 1.129 | 2008 |  |  |  |  | 45 |
| ${ }^{23} \mathrm{Na}_{2}$ | ${ }^{L^{1} \Sigma_{g_{+}}+}$ | 0 | 3.0405 | 163 | 0 | 0 | 5.12 | 0.73 |  |
|  | ${ }^{1} \Sigma^{\text {g }}{ }_{\text {g }}^{+}$ | 0 | 3.0789 | 159 |  |  | 4.90 | 0.72 | 1 |
| ${ }^{23} \mathrm{Na}_{2}{ }^{+}$ | $2 \Sigma^{g_{+}}$ 2 2 | 0 | 3.5725 | 107 | 0 | 0 |  | 1.04 0.96 |  |
| ${ }^{23} \mathrm{Na}_{2}{ }^{-}$ | ${ }_{2}^{2} \Sigma_{\text {a }}{ }^{\text {g }}{ }^{+}$ | 0 | 3.4723 | 99 | 0 | 0 | 0.46 | 0.96 0.64 | 1 |
|  |  |  |  |  |  |  | 0.43 | (0.44) | 1,19 |
| ${ }^{23} \mathrm{Na}{ }^{35} \mathrm{Cl}$ | ${ }^{1} \Sigma^{+}$ | 0 | 2.3758 | 357 | 45 | -8.75 | 9.07 | 3.97 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 2.3608 | 366 |  | 9.00 | 8.93 | $4.23 \pm 0.1$ | 1,19 |
| ${ }^{23} \mathrm{Na}^{35} \mathrm{Cl}^{+}$ | ${ }^{2} \Pi$ | 0 | 2.7970 | 157 | 32 | -6.76 |  | $0.33\left(\mathrm{Na}^{+}+\mathrm{Cl}\right)$ |  |
|  |  |  |  |  |  |  |  | 0.33 | 1 |
| ${ }^{23} \mathrm{Na}^{35} \mathrm{Cl}^{-}$ | ${ }^{2} \Sigma^{+}$ | 0 | 2.5158 | 263 | 44 | +4.03 | 0.85 | $1.31\left(\mathrm{Na}+\mathrm{Cl}^{-}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 2.497 | 265 |  |  | 0.73 | 1.34 | 40 |
| ${ }^{23} \mathrm{Na}^{19} \mathrm{~F}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.9407 | 526 | 56 | -8.00 |  | 4.71 |  |

## (Continued on next page)

Table 3 (Continued)

(Continued on next page)

Table 3 (Continued)

| Molecule | State | $T_{\text {e }}$ | $r_{\text {e }}$ | $\omega_{\mathrm{e}}$ | A | $\mu$ | $E_{\mathrm{i}}$ | $D_{0}$ | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{16} \mathrm{O}^{1} \mathrm{H}^{+}$ | ${ }^{2} \Pi$ | 0 | 0.9697 | 3738 | 265 | 1.67 | 13.01 | 4.39 | 1,30 |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.0372 | 3040 |  | +2.03 |  | $4.94\left(\mathrm{O}+\mathrm{H}^{+}\right)$ |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.0289 | 3113 |  | 2.32 |  | 5.09 | 1,22 |
| ${ }^{16} \mathrm{O}^{1} \mathrm{H}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 0.9666 | 3668 | 195 | +1.39 | 1.85 | $4.64\left(\mathrm{O}^{-}+\mathrm{H}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 0.964 | 3738 |  | 1.04 | 1.83 | 4.76 | 1,22,34 |
| ${ }^{31} \mathrm{P}_{2}$ | ${ }^{1} \Sigma_{\mathrm{g}}{ }^{+}$ | 0 | 1.8905 | 807 | 0 | 0 | 10.32 | 4.94 |  |
|  | ${ }^{1} \Sigma_{\text {g }}{ }^{+}$ | 0 | 1.8934 | 781 |  |  | 10.53 | 5.03 | 1 |
| ${ }^{31} \mathbf{P}_{\mathbf{2}}{ }^{+}$ | ${ }^{2} \Sigma_{\mathrm{g}}{ }^{\text {+ }}$ | 3408 | 1.8881 | 768 | 0 | 0 |  |  |  |
|  | ${ }^{2} \bar{\Pi}_{u}$ | 0 | 1.9787 | 701 | 0 | 0 |  | 5.01 |  |
|  | ${ }^{2} \Sigma_{\mathrm{g}}{ }^{+}$ | (2179) | 1.893 | $733\left(\omega_{0}\right)$ |  |  |  |  | 1 |
|  | ${ }^{2} \Pi_{u}$ | 0 | 1.9859 | 672 |  |  |  | 4.99 | 1 |
| ${ }^{31} \mathrm{P}_{\mathbf{2}}{ }^{-}$ | ${ }^{2} \Pi_{u}$ | 0 | 1.9830 | 674 | 0 | 0 | 0.60 | 4.82 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.979 | 640 |  |  | 0.59 |  | 3,34 |
| ${ }^{31} \mathrm{P}^{35} \mathrm{Cl}$ | ${ }^{3} \Sigma^{-}$ | 0 | 2.0393 | 528 | 64 | -0.59 |  | 3.38 |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 2.0146 | 551 |  |  |  | 2.96 | 19,50 |
| ${ }^{31} \mathbf{P}{ }^{19} \mathbf{F}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.6062 | 830 | 120 | -0.89 | 9.64 | 4.78 |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.5897 | 847 |  |  |  | $4.51 \pm 1.0$ | 1,19 |
| ${ }^{31} \mathbf{P}^{19} \mathbf{F}^{+}$ | ${ }^{2} \Pi$ | 0 | 1.5182 | 1022 | 115 | -2.61 |  | $5.52\left(\mathrm{P}^{+}+\mathrm{F}\right)$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.5003 | 1053 |  |  |  |  | 1 |
| ${ }^{31} \mathbf{P}^{1} \mathbf{H}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.4278 | 2338 | 76 | +0.42 | 10.17 | 3.19 |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.4214 | 2367 |  |  | 10.15 | 3.02 | 20,30,51 |
| ${ }^{31} \mathbf{P}^{1} \mathbf{H}^{+}$ | ${ }^{2} \Pi$ | 0 | 1.4331 | 2346 | 0 | +0.64 |  | 3.41 ( $\left.\mathrm{P}^{+}+\mathrm{H}\right)$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.4352 | $2300\left(\omega_{0}\right)$ |  |  |  | $\leq 3.36$ | 1 |
| ${ }^{31} \mathbf{P}^{1} \mathbf{H}^{-}$ | ${ }^{2} \Pi$ | 0 | 1.4420 | 2213 | 296 | +0.30 | 1.11 | $3.39\left(\mathrm{P}+\mathrm{H}^{-}\right)$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.407 | $2230\left(\omega_{0}\right)$ |  |  | 1.03 | 3.33 | 1,27 |
| ${ }^{31} \mathbf{P}^{14} \mathrm{~N}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.4828 | 1406 | 1 | -2.86 | 11.84 | 6.21 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.4909 | 1337 |  | 2.75 | 11.84 | $6.36 \pm 0.2$ | 1,19 |
| ${ }^{31} \mathbf{P}^{14} \mathrm{~N}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.4756 | 1368 | 8 | -2.05 |  | 4.75( $\left.\mathrm{P}^{+}+\mathrm{N}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 |  | (1200) |  |  |  |  | 1 |
| ${ }^{31} \mathbf{P}^{16} \mathbf{O}$ | ${ }^{2} \Pi$ | 0 | 1.4801 | 1258 | 48 | $\begin{gathered} -2.06 \\ 1.88 \end{gathered}$ | $\begin{gathered} 8.56 \\ 8.5 \end{gathered}$ | 6.13 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.4764 | 1233 |  |  |  | 6.15 | 1,43,52 |
| ${ }^{31} \mathbf{P}^{16} \mathbf{O}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.4249 | 1455 | 17 | -3.46 |  | $7.95\left(\mathrm{P}^{+}+\mathrm{O}\right)$$(8.4)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.4250 | 1412 |  |  |  |  | 1,43 |
| ${ }^{31} \mathrm{P}^{16} \mathrm{O}^{-}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.5499 | 1032 | 77 | -0.39 | 1.02 | $5.45\left(\mathrm{P}+\mathrm{O}^{-}\right)$ |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.540 | $1000\left(\omega_{0}\right)$ |  |  | 1.09 | 5.78 | 1 |
| ${ }^{31} \mathbf{P}^{32} \mathbf{S}$ | ${ }^{2} \Pi$ | 0 | 1.9056 | 745 | 20 | -0.68 | 8.06 | 4.44 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.8977 | 739 |  |  |  | 4.54 | 1,53 |
| ${ }^{31} \mathbf{P}^{32} \mathbf{S}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.8282 | 866 | 15 | -1.35 |  | $6.77\left(\mathrm{P}^{+}+\mathrm{S}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 |  | 845 |  |  |  |  | 1 |
| ${ }^{32} \mathrm{~S}_{2}$ | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{-}$ | 0 | 1.9070 | 715 | 0 | 0 | $\begin{aligned} & 9.55 \\ & 9.36 \end{aligned}$ | 4.39 |  |
|  | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{\text {- }}$ | 0 | 1.8892 | 726 |  |  |  | 4.37 | 1 |
| ${ }^{32} \mathrm{~S}_{2}{ }^{+}$ |  | 0 | 1.8269 | 827 | 0 | 0 |  | 5.385.37 |  |
|  |  | 0 | 1.825 | 790 |  |  |  |  | 1 |
| ${ }^{32} \mathrm{~S}_{2}{ }^{-}$ | ${ }_{2}^{2} \Pi_{\mathrm{g}}^{8}$ | 0 | 2.0284 | 564 | 0 | 0 | $\begin{aligned} & 1.57 \\ & 1.66 \end{aligned}$ | 3.94 |  |
|  | ${ }^{2} \Pi_{g}{ }^{2} \square^{\text {I }}$ | 0 |  | 601 558 |  |  |  | 2.852.83 | 1 |
| ${ }^{32} \mathrm{~S}{ }^{35} \mathrm{Cl}$ | ${ }^{2} \Pi$ |  | $\begin{aligned} & 2.0000 \\ & 1.975 \end{aligned}$ | 1.975 | 22 | -0.04 |  |  | 19,43 |
| ${ }^{32} \mathbf{S}^{19} \mathrm{~F}$ | ${ }^{2} \Pi$ | 0 | 1.6125 | 828 | 81 | $\begin{gathered} -0.85 \\ 0.87 \end{gathered}$ | $\begin{aligned} & 10.20 \\ & 10.09 \end{aligned}$ | 3.643.51 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.596 | 838 |  |  |  |  | 1,19,43 |
| ${ }^{32} \mathbf{S}^{19} \mathbf{F}^{+}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.5183 | 1001 | 82 | -2.18 |  | $3.99\left(\mathrm{~S}^{+}+\mathrm{F}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 19 |
| ${ }^{32} \mathbf{S}^{19} \mathbf{F}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.7413 | 629 | 64 | $+0.77$ | 2.13 | $2.90\left(\mathrm{~S}+\mathrm{F}^{-}\right)$ |  |
|  |  |  |  |  |  |  | 2.2810.46 |  | 19 |
| ${ }^{32} \mathbf{S}^{1} \mathbf{H}$ | ${ }^{2} \Pi$ | 0 | 1.3471 | 2664 | 1 | $\begin{gathered} +0.80 \\ 0.76 \end{gathered}$ |  | $\begin{aligned} & 3.65 \\ & 3.55 \end{aligned}$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.3409 | 2712 |  |  | 10.37 |  | 1,20,30 |

(Continued on next page)

Table 3 (Continued)

| Molecule | State | $T_{\text {e }}$ | $r_{\text {e }}$ | $\omega_{\text {e }}$ | A | $\mu$ | $E_{\text {i }}$ | $D_{0}$ | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{32} \mathbf{S}^{1} \mathbf{H}^{+}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.3714 | 2510 | 75 | +1.10 |  | 3.74( $\left.\mathrm{S}^{+}+\mathrm{H}\right)$ |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | $1.3744\left(r_{0}\right)$ |  |  |  |  | 3.48 | 1 |
| ${ }^{32} \mathbf{S}^{1} \mathbf{H}^{-}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.3485 | 2610 | 109 | $+0.60$ | 2.34 | $3.79\left(\mathrm{~S}^{-}+\mathrm{H}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.343 | 2647 |  |  | 2.32 | 3.79 | 1,34 |
| ${ }^{28} \mathrm{Si}_{2}$ | ${ }^{3} \Pi_{u}$ | 1215 | 2.1591 | 546 | 0 | 0 |  |  |  |
|  | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{-}$ | 0 | 2.2678 | 500 | 0 | 0 | 7.87 | 3.19 |  |
|  | ${ }^{3} \Pi_{\text {¢ }}$ |  | 2.155 | 548 |  |  |  |  | 1 |
|  | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{\text {u }}$ | 0 | 2.246 | 511 |  |  | 7.92 | 3.21 | 1,54 |
| ${ }^{28} \mathrm{Si}_{2}{ }^{+}$ | ${ }^{4} \Sigma_{\mathrm{g}}^{\mathrm{g}}$ | 0 | 2.2928 | 446 | 0 | 0 |  | 3.44 |  |
|  |  |  |  |  |  |  |  |  | 54 |
| ${ }^{28} \mathrm{Si}_{2}{ }^{-}$ | ${ }^{2} \Sigma_{\mathrm{g}}{ }^{+}$ | 0 | 2.1055 | 589 | 0 | 0 | 1.87 | 3.87 |  |
|  |  |  |  |  |  |  | 2.20 |  | 19 |
| ${ }^{28} \mathrm{Si}^{\mathbf{1 2}} \mathrm{C}$ | ${ }^{3} \Sigma^{-}$ | 2471 | 1.8078 | 868 | 0 | -2.58 |  |  |  |
|  | ${ }^{3} \Pi$ | 0 | 1.7138 | 982 | 17 | -1.73 | 8.89 | 4.27 |  |
|  | ${ }^{3} \Sigma^{-}$ | (5597) | (1.90) | (606) |  |  |  |  | 1 |
|  | ${ }^{3} \Pi$ | 0 | 1.722 | 965 |  |  | 9.0 | (4.64) | 1,55 |
| ${ }^{28} \mathrm{Si}^{\mathbf{1 2}} \mathrm{C}^{+}$ | ${ }^{4} \Sigma^{-}$ | 0 | 1.8257 | 812 | 29 | -1.26 |  | $3.50\left(\mathrm{Si}^{+}+\mathrm{C}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 55 |
| ${ }^{28} \mathrm{Si}^{35} \mathrm{Cl}$ | ${ }^{2} \Pi$ | 0 | 2.0880 | 507 | 104 | -0.96 | 7.40 | 4.30 |  |
|  | ${ }^{2} \Pi$ | 0 | 2.058 | 536 |  |  | 7.33 | 4.17 | 1,19,54,56 |
| ${ }^{28} \mathrm{Si}^{35} \mathrm{Cl}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.9677 | 650 | 127 | -3.35 |  | $5.02\left(\mathrm{Si}^{+}+\mathrm{Cl}\right)$ |  |
|  |  |  |  |  |  |  |  |  | 54 |
| ${ }^{28} \mathrm{Si}^{35} \mathrm{Cl}^{-}$ | ${ }^{3} \Sigma^{-}$ | 0 | 2.2761 | 342 | 70 | +2.17 | 1.19 | $1.98\left(\mathrm{Si}+\mathrm{Cl}^{-}\right)$ |  |
|  |  |  |  |  |  |  |  | $1.40 \pm 0.3$ | 32 |
| ${ }^{28} \mathbf{S i}{ }^{19} \mathrm{~F}$ | ${ }^{2} \Pi$ | 0 | 1.6196 | 833 | 127 | -0.96 | 7.47 | 6.07 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.6011 | 857 |  |  | 7.28 | 5.98 | 1,57 |
| ${ }^{28} \mathrm{Si}^{19} \mathrm{~F}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.5437 | 1016 | 131 | -3.43 |  | $6.71\left(\mathrm{Si}^{+}+\mathrm{F}\right)$ |  |
|  |  |  | 1.527 |  |  |  |  |  | 1,32 |
| ${ }^{28} \mathrm{Si}^{19} \mathrm{~F}^{-}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.7058 | 651 | 97 | +1.95 | 0.72 | $3.23\left(\mathrm{Si}+\mathrm{F}^{-}\right)$ |  |
|  |  |  |  |  |  |  | $1.34 \pm 0.3$ |  | 32 |
| ${ }^{28} \mathbf{S i}{ }^{1} \mathrm{H}$ | ${ }^{2} \Pi$ | 0 | 1.5275 | 2013 | 280 | $+0.07$ | 8.01 | 3.12 |  |
|  | ${ }^{2} \Pi$ | 0 | 1.5197 | 2042 |  |  | $\leq 8.04$ | $\leq 3.06$ | 1,43 |
| ${ }^{28} \mathrm{Si}^{1} \mathrm{H}^{+}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.5144 | 2126 | 76 | +0.17 |  | $3.22\left(\mathrm{Si}^{+}+\mathrm{H}\right)$ |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.5041 | 2157 |  |  |  | 3.17 | 1 |
| ${ }^{28} \mathrm{Si}^{1} \mathrm{H}^{-}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.5573 | 1816 | 559 | $+0.13$ | 1.28 | $3.05\left(\mathrm{Si}^{-}+\mathrm{H}\right)$ |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.5658 | 1804 |  |  | 1.28 | $\leq 2.95$ | 1,27 |
| ${ }^{28} \mathbf{S i}^{\mathbf{1 4}} \mathbf{N}$ | ${ }^{2} \Pi$ | 849 | 1.6400 | 1063 | 17 | -2.81 |  |  |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.5674 | 1183 | 25 | -2.78 |  | 4.61 |  |
|  | ${ }^{2} \Pi$ | 2032 | 1.642 | 1032 |  |  |  |  |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | 1.5719 | 1151 |  |  |  | 4.68 | 1,58 |
| ${ }^{28} \mathrm{Si}^{\mathbf{1 6}} \mathrm{O}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.5131 | 1258 | 50 | -3.13 | 11.45 | 8.04 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.5097 | 1242 |  | 3.10 | 11.49 | 8.26 | 1,19 |
| ${ }^{28} \mathrm{Si}^{16} \mathrm{O}^{+}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.5233 | 1143 | 2 | -4.39 |  | $4.71\left(\mathrm{Si}^{+}+\mathrm{O}\right)$ |  |
|  | ${ }^{2} \Sigma^{+}$ | 0 | $1.5191\left(r_{0}\right)$ |  |  |  |  | 4.98 | 1 |
| ${ }^{28} \mathrm{Si}^{31} \mathrm{P}$ | ${ }^{2} \Sigma^{+}$ | 0 | 1.9829 | 686 | 3 | -1.47 |  | 3.54 |  |
|  |  |  |  |  |  |  |  | 3.73 | 1 |
| ${ }^{28} \mathrm{Si}^{\mathbf{3 2}} \mathrm{S}$ | ${ }^{1} \Sigma^{+}$ | 0 | 1.9392 | 744 | 51 | -1.81 |  | 6.09 |  |
|  | ${ }^{1} \Sigma^{+}$ | 0 | 1.9293 | 750 |  | -1.73 |  | 6.42 | 1 |
| ${ }^{32} \mathrm{~S}{ }^{16} \mathrm{O}$ | ${ }^{3} \Sigma^{-}$ | 0 | 1.4919 | 1152 | 15 | -1.56 | 10.54 | 5.36 |  |
|  | ${ }^{3} \Sigma^{-}$ | 0 | 1.4811 | 1149 |  | -1.55 | 10.29 | 5.36 | 1 |
| ${ }^{32} \mathrm{~S}^{\mathbf{1 6}} \mathrm{O}^{+}$ | ${ }^{2} \Pi$ | 0 | 1.4240 | 1368 | 5 | -2.29 |  | 5.37( $\left.\mathrm{S}^{+}+\mathrm{O}\right)$ |  |
|  | ${ }^{2} \Pi$ | 0 | 1.424 | 1360 |  |  |  | 5.43 | 1 |
| ${ }^{32} \mathrm{~S}^{\mathbf{1 6}} \mathrm{O}^{-}$ | ${ }^{2} \Pi$ | 0 | 1.5901 | 888 | 26 | $-0.72$ | 0.96 | $4.31\left(\mathrm{~S}^{-}+\mathrm{O}\right)$ |  |
|  |  |  |  |  |  |  | 1.13 |  | 3 |

## 7. ISOTROPIC HYPERFINE COUPLING CONSTANTS OF POLYATOMIC MOLECULES

For this sensitive property, the contractions of the cc-pvqz basis sets were completely resolved.
For some diatomic molecules $\left(\mathrm{Cl}_{2}^{-}, \mathrm{ClF}^{-}, \mathrm{F}_{2}^{-}\right)$the calculations were performed at the calculated bond length (above) and the still uncertain experimental bond length (below in parentheses). Tables for calculated or experimental hfc constants are rare in review articles. Table 4 represents, to our knowledge, the largest set of hfc constants where 20 molecules were investigated to compare the hfc constants of 34 nuclei with experimental data. Irrespective of the meaning of the overall statistics, the mean absolute deviation $\backslash \overline{\Delta A_{i s o} \mid}$ resulted to be 10.1 G for an absolute range of $A_{\text {iso }}$ of 406 G . There are only two nuclei, Al in AlO and F in $\mathrm{ClF}^{-}$, for which the calculated values deviate markedly ( $\sim 100 \mathrm{G}$ ) from the experiment. When these two nuclei are removed from the statistics, the mean absolute deviation for $A_{\text {iso }}$ is reduced to 3.6 G . In particular, the nuclei $\mathrm{H}, \mathrm{B}, \mathrm{C}, \mathrm{N}$, and O exhibit good agreement between the calculated and experimental hfc constants.

Table 4 Calculated and experimental isotropic hyperfine coupling constants $A_{\text {iso }} / \mathrm{G}$.


Table 4 (Continued)

| Radical | State | Nucleus | $A_{\text {iso }}$ (calc.) | $A_{\text {iso }}($ exp.) | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | -26.2 |  |
| SH | ${ }^{2} \Pi$ | ${ }^{1} \mathrm{H}$ | -13.6 | -23.2 | 64 |
| BO | ${ }^{2} \Sigma^{+}$ | ${ }^{11} \mathrm{~B}$ | 381.3 | 365.7 | 65 |
| $\mathrm{O}_{2}$ | ${ }^{3} \Sigma_{\mathrm{g}}{ }^{-}$ | ${ }^{17} \mathrm{O}$ | -13.9 | -19.6 | 66 |
| $\mathrm{CH}_{3}$ | ${ }^{2} \mathrm{~A}_{2}{ }^{\prime}$ | $\begin{aligned} & { }^{13} \mathrm{C} \\ & { }^{1} \mathrm{H} \end{aligned}$ | $\begin{array}{r} 29.2 \\ -21.8 \end{array}$ | $\begin{gathered} 27 \\ -25 \end{gathered}$ | 67 |
| $\mathrm{NH}_{2}$ | ${ }^{2} \mathrm{~B}_{1}$ | $\begin{aligned} & { }^{14} \mathrm{~N} \\ & { }^{1} \mathrm{H} \end{aligned}$ | $\begin{array}{r} 9.9 \\ -20.4 \end{array}$ | $\begin{gathered} 9.96 \\ -24.0 \end{gathered}$ | 63 |
| $\mathrm{H}_{2} \mathrm{O}^{+}$ | ${ }^{2} \mathrm{~B}_{1}$ | $\begin{gathered} { }^{17} \mathrm{O} \\ { }^{1} \mathrm{H} \end{gathered}$ | $\begin{aligned} & -23.2 \\ & -23.7 \end{aligned}$ | $\begin{aligned} & -29.7 \\ & -26.1 \end{aligned}$ | 68 |
| $\mathrm{C}_{2} \mathrm{H}$ | ${ }^{2} \Sigma^{+}$ | $\begin{gathered} { }^{13} \mathrm{C} \\ { }^{13} \mathrm{C}(\mathrm{H}) \\ { }^{1} \mathrm{H} \end{gathered}$ | $\begin{array}{r} 375.6 \\ 80.7 \\ 18.8 \end{array}$ | $\begin{gathered} 362 \\ 76 \\ 18 \end{gathered}$ | 67 |
| $\mathrm{H}_{2} \mathrm{CN}$ | ${ }^{2} \mathrm{~B}_{1}$ | $\begin{gathered} { }^{13} \mathrm{C} \\ { }^{14} \mathrm{~N} \\ { }^{1} \mathrm{H} \end{gathered}$ | $\begin{array}{r} -24.4 \\ 8.1 \\ 81.3 \end{array}$ | $\begin{gathered} -29 \\ 9 \\ 83 \end{gathered}$ | 69 |
| $\mathrm{C}_{2} \mathrm{O}$ | ${ }^{3} \Sigma^{-}$ | $\begin{gathered} { }^{13} \mathrm{C} \\ { }^{13} \mathrm{C}(\mathrm{O}) \end{gathered}$ | $\begin{array}{r} 14.2 \\ -6.9 \end{array}$ | $\begin{gathered} 11.4 \\ -10.0 \end{gathered}$ | 70 |
| $\mathrm{C}_{3} \mathrm{H}_{5}$ (Allyl) | ${ }^{2} \mathrm{~A}_{2}$ | $\begin{gathered} { }^{13} \mathrm{C}\left(\mathrm{CH}_{2}\right) \\ { }^{13} \mathrm{C} \\ { }^{1} \mathrm{H}(\text { trans } \mathrm{HCCC}) \\ { }^{1} \mathrm{H} \text { (cis HCCC) } \\ { }^{1} \mathrm{H}(\mathrm{CH}) \end{gathered}$ | $\begin{array}{r} 18.6 \\ -16.0 \\ -14.4 \\ -13.5 \\ 4.2 \end{array}$ | $\begin{gathered} 21.9 \\ -17.2 \\ -14.8 \\ -13.9 \\ 4.2 \end{gathered}$ | 67 |
| $\mathrm{C}_{6} \mathrm{H}_{5} \mathrm{CH}_{2}$ | ${ }^{2} \mathrm{~B}_{1}$ | $\begin{gathered} { }^{13} \mathrm{C}\left(\mathrm{CH}_{2}\right) \\ { }^{13} \mathrm{C} \\ { }^{13} \mathrm{C}(\mathrm{o}) \\ { }^{13} \mathrm{C}(\mathrm{~m}) \\ { }^{13} \mathrm{C}(\mathrm{p}) \\ { }^{1} \mathrm{H}\left(\mathrm{CH}_{2}\right) \\ { }^{1} \mathrm{H}(\mathrm{o}) \\ { }^{1} \mathrm{H}(\mathrm{~m}) \\ { }^{1} \mathrm{H}(\mathrm{p}) \end{gathered}$ | $\begin{array}{r} 20.9 \\ -13.7 \\ 7.9 \\ -7.0 \\ 8.3 \\ -15.75 \\ -5.24 \\ 2.22 \\ -6.15 \end{array}$ | $\begin{gathered} -16.35 \\ -5.14 \\ 1.75 \\ -6.14 \end{gathered}$ | 71 |

${ }^{a}$ Uncertain experimental bond length

## 8. STRUCTURES OF POLYATOMIC GAS-PHASE MOLECULES

Experimentally obtained molecular structures (MW, ED) can be as uncertain as calculated structures. Even for the seemingly well-understood methane molecule $\mathrm{CH}_{4}$, the spectroscopically estimated CH
distance had to be corrected repeatedly over the years: 1966: $1.0940 \AA$ [72], 1985: $1.092 \AA$ [12], 1994: $1.086 \pm 0.002 \AA[73]$, and 1999: $1.0870 \AA$ [19]. For the ammonia-borane molecule $\mathrm{NH}_{3} \mathrm{BH}_{3}$ only two different values $r(\mathrm{BN})=1.56$ and $1.60 \AA$ from X-ray diffraction studies of the crystal structure [74] were available over almost 30 years for comparison with calculated structures [12]. In 1983, microwave spectroscopy suggested $r(\mathrm{BN})=1.6576 \AA$ [75] in excellent agreement with calculated structures. In Table 5, the calculated geometries of 53 polyatomic molecules are compared with experimental data; 74 out of a set of 109 calculated bond lengths differ by the order of $10^{-3} \AA$ or less from experimental values.

The worst cases are those compounds where sulfur and/or chlorine are involved. The calculated SF bond length in NSF is too long by $0.033 \AA$, compared with the experimental structure [12]; the calculated SCl bond length in $\mathrm{SCl}_{2}$ is too long by $0.037 \AA$, compared with the experimental structure [19]; and that in $\mathrm{OSCl}_{2}$ turned out to be too long by $0.048 \AA$, compared with the experimental structure [19]. The calculated ClO bond length in dichlorine peroxide $\mathrm{ClO}_{2} \mathrm{Cl}$ resulted to be too long by $0.052 \AA$, compared with the structure from microwave analysis [76]. For more details, see Section 11.

The $\mathrm{C}_{5} \mathrm{O}$ molecule is an example where two arguments favor the calculated structure, but not that from the recent microwave analysis [77]. At first, the sequence of the CC distances should be in accordance with the commonly accepted polyacetylene structure ${ }^{+}: \mathrm{O} \equiv \mathrm{C} 1-\mathrm{C} 2 \equiv \mathrm{C} 3-\mathrm{C} 4 \equiv \mathrm{C} 5:^{-}$; and second, the calculated rotational constant $(1369 \mathrm{MHz})$ is in almost perfect agreement with the observed one (1367 MHz ).

Another example where an experimental bond length is rebutted in the literature is $\mathrm{NH}_{3} \mathrm{BF}_{3}$. Similar to $\mathrm{NH}_{3} \mathrm{BH}_{3}$, the BN bond length in $\mathrm{NH}_{3} \mathrm{BF}_{3}$ is likely $0.06-0.10 \AA$ larger than the hitherto assumed value of $r(\mathrm{BN})=1.59 \AA[12,78]$ in agreement with calculated values.

A difficult geometry parameter is the $\mathrm{SO}_{\mathrm{H}}$ bond length in the sulfuric acid molecule $\mathrm{H}_{2} \mathrm{SO}_{4}$. Calculated values from the literature overestimate that from microwave spectroscopy by $0.048 \AA$. B3LYP density functional theory in conjunction with extended basis sets (cc-pvqz) reduces this discrepancy to $0.026 \AA$ (cf. Table 5).

Table 5 Calculated and experimental equilibrium geometries (interatomic distances $r / \AA$ and angles </ degrees) for selected polyatomic gas-phase molecules.


Table 5 (Continued)

| Molecule | Point Group | Geometrical Parameter | Calc. | Exp. | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $r(\mathrm{CH})$ | 1.0608 | 1.060 |  |
| $\mathrm{C}_{2} \mathrm{H}_{4}$ | $\mathrm{D}_{2 \mathrm{~h}}$ | $\begin{gathered} r(\mathrm{CC}) \\ r(\mathrm{CH}) \\ <(\mathrm{HCH}) \end{gathered}$ | $\begin{gathered} 1.3240 \\ 1.0819 \\ 116.5 \end{gathered}$ | $\begin{aligned} & 1.339 \\ & 1.087 \\ & 117.4 \end{aligned}$ | 19 |
| $\mathrm{C}_{2} \mathrm{H}_{6}$ | $\mathrm{D}_{3 \mathrm{~d}}$ |  | $\begin{gathered} 1.5269 \\ 1.0904 \\ 107.5 \end{gathered}$ | $\begin{gathered} 1.5351 \\ 1.0940 \\ 107.8 \end{gathered}$ | 19 |
| HCN | $\mathrm{C}_{\infty \mathrm{V}}$ | $\begin{aligned} & r(\mathrm{CN}) \\ & r(\mathrm{CH}) \end{aligned}$ | $\begin{aligned} & 1.1450 \\ & 1.0654 \end{aligned}$ | $\begin{aligned} & 1.1532 \\ & 1.0655 \end{aligned}$ | 19 |
| HCP | $\mathrm{C}_{\infty \mathrm{v}}$ | $\begin{aligned} & r(\mathrm{CP}) \\ & r(\mathrm{CH}) \end{aligned}$ | $\begin{aligned} & 1.5338 \\ & 1.0700 \end{aligned}$ | $\begin{aligned} & 1.5398 \\ & 1.0692 \end{aligned}$ | 19 |
| $\mathrm{CH}_{3} \mathrm{OH}$ | $\mathrm{C}_{\mathrm{s}}$ | $\begin{gathered} r(\mathrm{CO}) \\ r\left(\mathrm{CH}_{\mathrm{tr}}\right) \\ r\left(\mathrm{CH}_{\mathrm{g}}\right) \\ r(\mathrm{OH}) \\ <\left(\mathrm{OCH}_{\mathrm{tr}}\right) \\ <\left(\mathrm{OCH}_{\mathrm{g}, \mathrm{~g}^{\prime}}\right)^{\mathrm{a}} \\ <\left(\mathrm{H}_{\mathrm{g}} \mathrm{CH}_{\mathrm{g}^{\prime}}\right) \\ <(\mathrm{COH}) \end{gathered}$ | $\begin{gathered} 1.4201 \\ 1.0878 \\ 1.0939 \\ 0.9594 \\ 106.9 \\ 130.6 \\ 108.9 \\ 108.9 \end{gathered}$ | 1.4246 1.0936 (av.) $"$ 0.9451 107.2 129.9 108.63 108.53 | 12,19 |
| $\mathrm{CH}_{3} \mathrm{SH}$ | $\mathrm{C}_{\mathrm{s}}$ | $\begin{gathered} r\left(\mathrm{CS}^{2}\right) \\ r\left(\mathrm{CH}_{\mathrm{tr}}\right) \\ r\left(\mathrm{CH}_{\mathrm{g}}\right) \\ r\left(\mathrm{SH}^{2}\right) \\ <\left(\mathrm{SCH}_{\mathrm{tr}}\right) \\ <\left(\mathrm{SCH}_{\mathrm{g}, \mathrm{~g}}\right)^{\mathrm{a}} \\ \left.<\left(\mathrm{H}_{\mathrm{g} H}^{\mathrm{g}}\right)^{\prime}\right) \\ <(\mathrm{CSH}) \end{gathered}$ | $\begin{gathered} 1.8261 \\ 1.0870 \\ 1.0862 \\ 1.3429 \\ 106.1 \\ 129.5 \\ 110.4 \\ 97.3 \end{gathered}$ | $\begin{gathered} 1.819 \\ 1.091 \\ 1.091 \\ 1.336 \\ - \\ - \\ 109.8 \\ 96.5 \end{gathered}$ | 12 |
| $\mathrm{CH}_{3} \mathrm{~F}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{gathered} r(\mathrm{CF}) \\ r(\mathrm{CH}) \\ <(\mathrm{HCH}) \end{gathered}$ | $\begin{gathered} 1.3883 \\ 1.0892 \\ 110.0 \end{gathered}$ | $\begin{gathered} 1.382 \\ 1.095 \\ 110.45 \end{gathered}$ | 19 |
| $\mathrm{NH}_{3} \mathrm{BH}_{3}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{gathered} r(\mathrm{BN}) \\ r(\mathrm{BH}) \\ r(\mathrm{NH}) \\ <(\mathrm{HBH}) \\ <(\mathrm{HNH}) \end{gathered}$ | $\begin{gathered} 1.6563 \\ 1.2063 \\ 1.0143 \\ 113.5 \\ 107.9 \end{gathered}$ | $\begin{gathered} 1.6576 \\ 1.2160 \\ 1.0140 \\ 113.8 \\ 108.6 \end{gathered}$ | 75 |
| $\mathrm{NH}_{2} \mathrm{OH}$ | $\begin{gathered} \mathrm{C}_{\mathrm{s}} \\ \text { trans } \end{gathered}$ | $\begin{gathered} r(\mathrm{NO}) \\ r(\mathrm{NH}) \\ r(\mathrm{OH}) \\ <(\mathrm{ONH} \\ <(\mathrm{an}, \mathrm{a})^{\mathrm{a}} \end{gathered}$ | $\begin{gathered} 1.4426 \\ 1.0155 \\ 0.9603 \\ 113.8 \\ 105.9 \end{gathered}$ | $\begin{aligned} & 1.453 \\ & 1.016 \\ & 0.962 \\ & 112.6 \\ & 107.1 \end{aligned}$ | 12 |
|  |  |  |  | (Continued on next page) |  |

Table 5 (Continued)


Table 5 (Continued)

| Molecule | Point <br> Group | Geometrical Parameter | Calc. | Exp. | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{HNO}_{2}$ | $\underset{\text { C }}{\mathrm{C}_{\mathrm{s}}} \text { trans }$ | $\begin{gathered} r(\mathrm{NO}) \\ r\left(\mathrm{NO}_{\mathrm{H}}\right) \\ r(\mathrm{OH}) \\ <(\mathrm{ONO}) \\ <(\mathrm{NOH}) \end{gathered}$ | $\begin{gathered} 1.1627 \\ 1.4297 \\ 0.9669 \\ 111.1 \\ 102.7 \end{gathered}$ | $\begin{aligned} & 1.170 \\ & 1.432 \\ & 0.958 \\ & 110.7 \\ & 102.1 \end{aligned}$ | 19 |
| $\mathrm{HNO}_{3}$ | $\mathrm{C}_{\mathrm{s}}$ planar | $\begin{gathered} r\left(\mathrm{NO}_{\text {cis }}\right) \\ r\left(\mathrm{NO}_{\text {tr }}\right) \\ r\left(\mathrm{NO}_{\mathrm{H}}\right) \\ r(\mathrm{OH}) \\ <\left(\mathrm{O}_{\text {cis }} \mathrm{NO}_{\mathrm{H}}\right) \\ <\left(\mathrm{O}_{\mathrm{tr}} \mathrm{NO}_{\mathrm{H}}\right) \\ <(\mathrm{NOH}) \end{gathered}$ | $\begin{gathered} 1.2066 \\ 1.1913 \\ 1.4097 \\ 0.9701 \\ 115.6 \\ 114.0 \\ 103.1 \end{gathered}$ | $\begin{gathered} 1.21 \\ 1.20 \\ 1.41 \\ 0.96 \\ 115.9 \\ 113.9 \\ 102.2 \end{gathered}$ | 19 |
| $\mathrm{FN}_{3}$ | $\underset{\text { trans }}{\mathrm{C}_{\mathrm{s}}}$ | $\begin{gathered} r(\text { FN1 }) \\ r(\text { N1N2 }) \\ r(\text { N2N3 }) \\ <(\text { N1N2N3 }) \\ <(\text { FN1N2 }) \end{gathered}$ | $\begin{gathered} 1.4248 \\ 1.2513 \\ 1.1267 \\ 171.8 \\ 105.3 \end{gathered}$ | $\begin{aligned} & 1.444 \\ & 1.253 \\ & 1.132 \\ & 170.9 \\ & 103.8 \end{aligned}$ | 81 |
| $\mathrm{N}_{2} \mathrm{~F}_{2}$ | $\mathrm{C}_{2 \mathrm{~h}}$ <br> trans | $\begin{gathered} r(\mathrm{NN}) \\ r(\mathrm{NF}) \\ <(\mathrm{NNF}) \end{gathered}$ | $\begin{gathered} 1.2146 \\ 1.3844 \\ 105.7 \end{gathered}$ | $\begin{aligned} & 1.214 \\ & 1.384 \\ & 114.5 \end{aligned}$ | 12 |
| $\mathrm{ClN}_{3}$ | $\begin{gathered} \mathrm{C}_{\mathrm{s}} \\ \text { trans } \end{gathered}$ | $\begin{gathered} r(\text { ClN1 }) \\ r(\text { (N1N2 }) \\ r(\text { N2N3 }) \\ <(\text { N1N2N3 }) \\ <(\text { ClN1N2 }) \end{gathered}$ | $\begin{gathered} 1.7610 \\ 1.2420 \\ 1.1265 \\ 172.1 \\ 110.1 \end{gathered}$ | $\begin{aligned} & 1.745 \\ & 1.252 \\ & 1.133 \\ & 171.9 \\ & 108.6 \end{aligned}$ | 81 |
| $\mathrm{C}_{4} \mathrm{H}_{6}$ | $\mathrm{C}_{2 \mathrm{~h}}$ | $\begin{gathered} r(\mathrm{C} 1 \mathrm{C} 2) \\ r(\mathrm{C} 2 \mathrm{C} 3) \\ r(\mathrm{C} 1 \mathrm{H}) \\ r\left(\mathrm{C} 1 \mathrm{H}^{\prime}\right) \\ r(\mathrm{C} 2 \mathrm{H}) \\ <(\mathrm{C} 1 \mathrm{C} 2 \mathrm{C} 3) \\ <\left(\mathrm{HC}^{\prime} 1 \mathrm{C} 2\right) \\ <\left(\mathrm{H}^{\prime} \mathrm{C} 1 \mathrm{C} 2\right) \\ <\left(\mathrm{HC}^{2} 1\right) \end{gathered}$ | $\begin{gathered} 1.3336 \\ 1.4526 \\ 1.0802 \\ 1.0826 \\ 1.0852 \\ 124.3 \\ 121.7 \\ 121.4 \\ 119.3 \end{gathered}$ | $\begin{gathered} 1.349 \\ 1.467 \\ 1.08 \text { (av.) } \\ " \\ " \\ 124.4 \\ 120.9 \text { (av.) } \\ " \\ " \end{gathered}$ | 19 |
| $\mathrm{C}_{3} \mathrm{O}$ | $\mathrm{C}_{\infty \mathrm{V}}$ | $\begin{gathered} r(\mathrm{C} 1 \mathrm{O}) \\ r(\mathrm{C} 1 \mathrm{C} 2) \\ r(\mathrm{C} 2 \mathrm{C} 3) \end{gathered}$ | $\begin{aligned} & 1.1488 \\ & 1.2942 \\ & 1.2670 \end{aligned}$ | $\begin{aligned} & 1.150 \\ & 1.306 \\ & 1.254 \end{aligned}$ | 82 |
| $\mathrm{C}_{5} \mathrm{O}$ | $\mathrm{C}_{\infty \mathrm{V}}$ | $\begin{gathered} r(\mathrm{C} 1 \mathrm{O}) \\ r(\mathrm{C} 1 \mathrm{C} 2) \\ r(\mathrm{C} 2 \mathrm{C} 3) \\ r(\mathrm{C} 3 \mathrm{C} 4) \end{gathered}$ | $\begin{aligned} & 1.1555 \\ & 1.2828 \\ & 1.2592 \\ & 1.2897 \end{aligned}$ | $\begin{aligned} & 1.1562 \\ & 1.2552 \\ & 1.2881 \\ & 1.2947 \end{aligned}$ | 77 |
|  |  |  |  | (Continued on next page) |  |

Table 5 (Continued)

| Molecule | Point <br> Group | Geometrical Parameter | Calc. | Exp. | Ref.(Exp.) |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $r$ (C4C5) | 1.2744 | 1.2736 |  |
| $\mathrm{OC}_{3} \mathrm{O}$ | $\mathrm{D}_{\infty}$ | $\begin{aligned} & r(\mathrm{CO}) \\ & r(\mathrm{CC}) \end{aligned}$ | $\begin{aligned} & 1.1597 \\ & 1.2715 \end{aligned}$ | $\begin{aligned} & 1.1632 \\ & 1.2894 \end{aligned}$ | 83 |
| $\mathrm{C}_{2} \mathrm{~S}$ | $\mathrm{C}_{\infty \mathrm{v}}$ | $\begin{aligned} & r(\mathrm{CS}) \\ & r(\mathrm{CC}) \end{aligned}$ | $\begin{aligned} & 1.5712 \\ & 1.3094 \end{aligned}$ | $\begin{aligned} & 1.567 \\ & 1.310 \end{aligned}$ | 84 |
| $\mathrm{C}_{3} \mathrm{~S}$ | $\mathrm{C}_{\text {ov }}$ |  | $\begin{aligned} & 1.5426 \\ & 1.2892 \\ & 1.2749 \end{aligned}$ | $\begin{aligned} & 1.5323 \\ & 1.3028 \\ & 1.2724 \end{aligned}$ | 85 |
| $\mathrm{SC}_{3} \mathrm{~S}$ | $\mathrm{D}_{\infty}$ | $\begin{aligned} & r(\mathrm{CS}) \\ & r(\mathrm{CC}) \end{aligned}$ | $\begin{aligned} & 1.5595 \\ & 1.2728 \end{aligned}$ | $\begin{aligned} & 1.56 \\ & 1.28 \end{aligned}$ | 86 |
| $\mathrm{ClO}_{2} \mathrm{Cl}$ | $\mathrm{C}_{2}$ | $\begin{gathered} r(\mathrm{OO}) \\ r(\mathrm{OCl}) \\ <(\mathrm{ClOO}) \\ <(\mathrm{ClOOCl}) \end{gathered}$ | $\begin{gathered} 1.3593 \\ 1.7559 \\ 111.5 \\ 84.9 \end{gathered}$ | $\begin{gathered} 1.4259 \\ 1.7044 \\ 110.07 \\ 81.03 \end{gathered}$ | 76 |
| $\mathrm{NO}_{3}{ }^{-}$ | $\mathrm{D}_{3 \mathrm{~h}}$ | $r$ (NO) | 1.2558 | $1.22-1.27$ | 12 |
| $\mathrm{BF}_{3}$ | $\mathrm{D}_{3 \mathrm{~h}}$ | $r$ (BF) | 1.3134 | 1.313 | 19 |
| $\mathrm{NF}_{3}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{gathered} r(\mathrm{NF}) \\ <(\mathrm{FNF}) \end{gathered}$ | $\begin{gathered} 1.3774 \\ 102.0 \end{gathered}$ | $\begin{aligned} & 1.365 \\ & 102.5 \end{aligned}$ | 12 |
| $\mathrm{PF}_{3}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{gathered} r(\mathrm{PF}) \\ <(\mathrm{FPF}) \end{gathered}$ | $\begin{gathered} 1.5831 \\ 97.6 \end{gathered}$ | $\begin{gathered} 1.570 \\ 97.8 \end{gathered}$ | 19 |
| $\mathrm{P}_{4}$ | $\mathrm{T}_{\mathrm{d}}$ | $r(\mathrm{PP})$ | 2.2077 | 2.21 | 19 |
| $\mathrm{NH}_{3} \mathrm{BF}_{3}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{gathered} r(\mathrm{BN}) \\ r(\mathrm{BF}) \\ r(\mathrm{NH}) \\ <(\mathrm{FBF}) \\ <(\mathrm{HNH}) \end{gathered}$ | $\begin{gathered} 1.6943 \\ 1.3703 \\ 1.0148 \\ 114.3 \\ 108.7 \end{gathered}$ | $\begin{gathered} 1.59(?)^{\mathrm{b}} \\ 1.38 \\ - \\ 111.0 \end{gathered}$ | 12,78 |
| $\mathrm{H}_{2} \mathrm{SO}_{4}$ | $\mathrm{C}_{2}$ | $\begin{gathered} r(\mathrm{SO}) \\ r\left(\mathrm{SO}_{\mathrm{H}}\right) \\ r(\mathrm{OH}) \\ <(\mathrm{OSO}) \\ <\left(\mathrm{O}_{\mathrm{H}} \mathrm{SO}_{\mathrm{H}}\right) \\ <\left(\mathrm{HOS}^{2}\right) \\ <\left(\mathrm{OOSO}_{\mathrm{H}}\right) \\ <\left(\mathrm{OSO}_{\mathrm{H}} \mathrm{H}\right) \end{gathered}$ | $\begin{gathered} 1.4257 \\ 1.5997 \\ 0.9674 \\ 124.0 \\ 102.0 \\ 109.0 \\ -125.1 \\ -29.7 \end{gathered}$ | $\begin{gathered} 1.422 \\ 1.574 \\ 0.97 \\ 123.3 \\ 101.3 \\ 108.5 \\ \\ -20.8 \end{gathered}$ | 19 |
| $\mathrm{ONF}_{3}$ | $\mathrm{C}_{3 \mathrm{v}}$ | $\begin{aligned} & r(\mathrm{NO}) \\ & r(\mathrm{NF}) \end{aligned}$ | $\begin{aligned} & 1.1516 \\ & 1.4385 \end{aligned}$ | $\begin{aligned} & 1.158 \\ & 1.431 \end{aligned}$ | $86$ <br> Continued on |

Table 5 (Continued)

| Molecule | Point <br> Group | Geometrical <br> Parameter | Calc. | Exp. | Ref.(Exp.) |
| :--- | :---: | :---: | :---: | :---: | :---: |
|  |  | $<(\mathrm{FNF})$ | 100.5 | 100.8 |  |
| $\mathrm{OSCl}_{2}$ | $\mathrm{C}_{\mathrm{s}}$ | $r(\mathrm{SO})$ | 1.4418 | 1.44 | 19 |
|  |  | $r(\mathrm{SCl})$ | 2.1202 | 2.072 |  |
|  |  | $<(\mathrm{OSCl})$ | 107.8 | 108.0 |  |
|  |  | $<(\mathrm{ClSCl})$ | 98.2 | 97.2 |  |
| $\mathrm{SF}_{6}$ | $\mathrm{O}_{\mathrm{h}}$ | $r(\mathrm{SF})$ | 1.5827 | 1.561 | 19 |

${ }^{a}$ The notation $<\left(\mathrm{ABH}_{d, \mathrm{e}}\right)$ is used to describe the angle between the AB bond and the $H_{d} B H_{e}$ plane.
${ }^{\mathrm{b}}$ Question-marks are taken from the literature.

## 9. ANHARMONICITIES OF DIATOMIC MOLECULES

For the calculation of anharmonicities, the potential energy curves were represented by means of a point grid of $0.1 a_{0}$ and roughly 25 points. These were interpolated by a polynomial fit from which the anharmonicities $\omega_{\mathrm{e}} x_{\mathrm{e}}$ within the framework of a Dunham analysis were derived. All experimental values are taken from ref. 1. The comparison of calculated and experimental values in Table 6 indicates a remarkably strong relation between anharmonicities and reduced masses of the vibrating systems and a surprisingly good agreement. Anharmonicities are useful for deriving $\omega_{0}$ values according to $\omega_{0}=\omega_{\mathrm{e}}-2 \omega_{\mathrm{e}} x_{\mathrm{e}}$. The statistics of deviations between calculated and experimental anharmonicities can be found in Table 2.

Table 6 Calculated and experimental anharmonicities $\omega_{\mathrm{e}} x_{\mathrm{e}} / \mathrm{cm}^{-1}$.

| Molecule | $\omega_{\mathrm{e}} x_{\mathrm{e}}$ (calc.) | $\omega_{\mathrm{e}} x_{\mathrm{e}}($ exp.) |
| :--- | :---: | :---: |
| ${ }^{27} \mathrm{Al}^{1} \mathrm{H}$ | 27 | 29.09 |
| ${ }^{9} \mathrm{Be}^{1} \mathrm{H}$ | 35 | 36.31 |
| ${ }^{9} \mathrm{Be}^{1} \mathrm{H}^{+}$ | 35 | 39.79 |
| ${ }^{9} \mathrm{Be}^{16} \mathrm{O}$ | 11 | 11.83 |
| ${ }^{11} \mathrm{~B}^{19} \mathrm{~F}$ | 12 | 11.84 |
| ${ }^{11} \mathrm{~B}^{1} \mathrm{H}$ | 47 | 49.39 |
| ${ }^{11} \mathrm{~B}^{16} \mathrm{O}$ | 11 | 11.81 |
| ${ }^{12} \mathrm{C}^{19} \mathrm{~F}$ | 11 | 11.10 |
| ${ }^{12} \mathrm{C}^{1} \mathrm{H}$ | 59 | 63.02 |
| ${ }^{12} \mathrm{C}^{14} \mathrm{~N}$ | 13 | 13.09 |
| ${ }^{12} \mathrm{C}^{14} \mathrm{~N}^{+}$ | 15 | 16.14 |
| ${ }^{12} \mathrm{C}^{16} \mathrm{O}$ | 13 | 13.29 |
| ${ }^{1} \mathrm{H}_{2}$ | 115 | 121.34 |
| ${ }^{1} \mathrm{H}^{35} \mathrm{Cl}$ | 50 | 52.82 |
| ${ }^{4} \mathrm{He}^{1} \mathrm{H}^{+}$ | 153 | 157.7 |
| ${ }^{1} \mathrm{H}^{19} \mathrm{~F}$ | 87 | 89.88 |
| ${ }^{7} \mathrm{Li}^{1} \mathrm{H}$ | 21 | 23.20 |
| ${ }^{24} \mathrm{Mg}^{1} \mathrm{H}$ | 32 | 31.89 |
| ${ }^{24} \mathrm{Mg}^{1} \mathrm{H}^{+}$ | 29 | 31.93 |
| ${ }^{14} \mathrm{~N}_{2}$ | 13 | 14.32 |
| ${ }^{14} \mathrm{~N}_{2}+$ | 14 | 16.10 |
| ${ }^{23} \mathrm{Na}^{1} \mathrm{H}$ | 16 | 19.72 |

(Continued on next page)

Table 6 (Continued)

| Molecule | $\omega_{\mathrm{e}} x_{\mathrm{e}}$ (calc.) | $\omega_{\mathrm{e}} x_{\mathrm{e}}$ (exp.) |
| :--- | :---: | :---: |
| ${ }^{14} \mathrm{~N}^{1} \mathrm{H}$ | 71 | 78.35 |
| ${ }^{14} \mathrm{~N}^{16} \mathrm{O}$ | 13 | 14.07 |
| ${ }^{14} \mathrm{~N}^{16} \mathrm{O}^{+}$ | 14 | 16.26 |
| ${ }^{16} \mathrm{O}_{2}$ | 10 | 11.98 |
| ${ }^{16} \mathrm{O}_{2}+$ | 13 | 16.26 |
| ${ }^{16} \mathrm{O}^{1} \mathrm{H}$ | 78 | 84.88 |
| ${ }^{16} \mathrm{O}^{1} \mathrm{H}^{+}$ | 81 | 78.52 |
| ${ }^{31} \mathrm{P}^{1} \mathrm{H}$ | 39 | 44.5 |
| ${ }^{28} \mathrm{Si}^{1} \mathrm{H}$ | 34 | 35.51 |
| ${ }^{32} \mathrm{~S}^{1} \mathrm{H}$ | 45 | $(59.9)$ |

## 10. NMR CHEMICAL SHIFTS OF POLYATOMIC MOLECULES

The calculation of NMR chemical shifts is confronted with "easy" and "difficult" (quadrupole effects) nuclei as well as with "normal" and "unusual" (diradicaloid character) bonding situations. Experience has shown that computational procedures such as GIAO (gauge-independent atomic orbitals) have much success where ${ }^{1} \mathrm{H}$ and ${ }^{13} \mathrm{C}$ nuclei are involved in the molecules (Tables 7a, 7b). Calculated chemical shifts of the nuclei ${ }^{15} \mathrm{~N},{ }^{17} \mathrm{O}$, and ${ }^{19} \mathrm{~F}$ are also in good agreement with experimental values for many compounds. However, there are exceptional molecules where all attempts of calculations fail dramatically. The same is true for ${ }^{29} \mathrm{Si},{ }^{31} \mathrm{P},{ }^{33} \mathrm{~S}$, and ${ }^{35} \mathrm{Cl}$. Worst cases are defined here by discrepancies of the size

Table 7a Relative shifts $\delta\left({ }^{1} \mathrm{H}\right)$ referenced to TMS; $\mathrm{CH}_{4}$ was used as primary reference for the calculated data, a value of $\delta\left(\mathrm{CH}_{4}\right)=0.1 \mathrm{ppm}$ was used for conversion, the calculated shielding $\sigma\left({ }^{1} \mathrm{H}\right)$ of $\mathrm{CH}_{4}$ is 31.5 ppm . Experimental values are for gas-phase molecules unless otherwise stated.

|  |  | $\delta\left({ }^{1} \mathrm{H}\right)($ calc. $)$ |
| :--- | :---: | :---: |
| $\mathrm{CH}_{4}$ | 0.1 | $\delta\left({ }^{1} \mathrm{H}\right)($ exp. $)$ |
| $\mathrm{H}_{2} \mathrm{O}$ | 0.5 | 0.1 |
| $\mathrm{NH}_{3}$ |  | -0.2 |
| HF | 2.3 | 0.7 |
| $\mathrm{PH}_{3}$ | 1.9 | 0.1 |
| $\mathrm{SiH}_{4}$ | 3.5 | 2.6 |
| $\mathrm{H}_{2} \mathrm{~S}$ | 0.8 | 1.6 |
| $\mathrm{HCl}^{2}$ | 0.5 | 3.1 |
| $\mathrm{C}_{2} \mathrm{H}_{2}$ |  | 1.0 |
| $\mathrm{C}_{2} \mathrm{H}_{4}$ | 5.6 | 0.2 |
| $\mathrm{C}_{2} \mathrm{H}_{6}$ |  | 0.8 |
| $\mathrm{HCN}^{\mathrm{CH}_{3} \mathrm{OH}}$ |  | 2.3 |
| $\mathrm{CH}_{3} \mathrm{~F}$ |  | 3.5 |
| $\mathrm{HOF}^{\mathrm{CH}} 3$ |  | -0.4 |
|  | 4.3 | 1.5 |
|  |  | 11.8 |

$\mid \delta($ calc. $)-\delta($ exp. $)|=|\Delta \delta|>100 \mathrm{ppm}$ and $100 \%| \Delta \delta| | \mid \delta($ exp. $) \mid>10 \%$.
All references for experimental NMR chemical shifts in the Tables 7a-h can be found in ref. 87. The relation between absolute shieldings $\sigma$ and relative shifts $\delta$ is $\delta=\sigma_{\text {ref }}-\sigma$ where all data are given in ppm . All reported relative NMR chemical shifts $\delta$ are isotropic shifts, $\delta_{\text {iso }}=1 / 3\left(\delta_{\mathrm{xx}}+\delta_{\mathrm{yy}}+\delta_{\mathrm{zz}}\right)$.

The calculated relative shift $\sigma\left({ }^{15} \mathrm{~N}\right)=1017 \mathrm{ppm}$ in the NSF molecule differs by $752 \mathrm{ppm}(284 \%)$ from the experimental value of 265 ppm , according to IGLO(II) ab initio RHF calculations from the literature. Recent GIAO/B3LYP/cc-pvqz calculations (Table 7c) reduce this error from 752 ( $284 \%$ ) to 455 $\mathrm{ppm}(172 \%)$. Difficult molecules for the calculation of relative shifts $\delta\left({ }^{17} \mathrm{O}\right)$ are $\mathrm{O}_{3}$, but also $\mathrm{SO}_{2}$ and $\mathrm{OSCl}_{2}$ (Table 7d). Different attempts at calculations for $\mathrm{O}_{3}$ are summarized in Table 7h. Ab initio RHF calculations for $\mathrm{O}_{3}$ fail by $\sim 10^{3} \mathrm{ppm}$. Such errors were attributed to the lack of electron correlation in

Table 7b Relative shifts $\delta\left({ }^{13} \mathrm{C}\right)$ referenced to TMS; $\mathrm{CH}_{4}$ was used as primary reference for the calculated data, a value of $\delta\left(\mathrm{CH}_{4}\right)=-7.0 \mathrm{ppm}$ was used for conversion, the calculated shielding $\sigma\left({ }^{13} \mathrm{C}\right)$ of $\mathrm{CH}_{4}$ is 189.0 ppm . Experimental values are for gas-phase molecules unless otherwise stated.

|  | $\delta\left({ }^{13} \mathrm{C}\right)$ (calc.) | $\delta\left({ }^{13} \mathrm{C}\right)$ (exp.) |
| :--- | :---: | :---: |
| $\mathrm{CH}_{4}$ | -7.0 | -7.0 |
| HCN | 109.7 | 106.0 |
| $\mathrm{CH}_{3} \mathrm{OH}$ | 55.2 | 51.5 |
| $\mathrm{CH}_{3} \mathrm{~F}$ | 75.6 | 71.3 |
| $\mathrm{C}_{2} \mathrm{H}_{2}$ | 73.0 | 70.9 |
| $\mathrm{C}_{2} \mathrm{H}_{4}$ | 131.4 | 123.6 |
| $\mathrm{C}_{2} \mathrm{H}_{6}$ | 9.1 | 7.2 |
| $\mathrm{CO}_{2}$ | 131.5 | 129.3 |
| $\mathrm{CS}_{2}$ | 212.8 | 196.1 |
| $\mathrm{CO}^{\mathrm{OCS}}$ |  | 197.4 |
| $\mathrm{CFCl}_{3}$ | 164.1 | 187.1 |
| $\left(\mathrm{CH}_{3}\right)_{3} \mathrm{C}^{+}$ |  | 157.8 |
| $\mathrm{H}_{2} \mathrm{CCCH}_{2}$ |  | 356.8 |
|  | 51.6 | 158.1 |
| $\mathrm{CH}_{3} \mathrm{SH}^{+}$ | C | 117.6 (liq.) |
|  |  | 229.4 |

Table 7c Relative shifts $\delta\left({ }^{15} \mathrm{~N}\right)$ referenced to $\mathrm{NH}_{3}$; the calculated shielding $\sigma\left({ }^{15} \mathrm{~N}\right)$ of $\mathrm{NH}_{3}$ is 262.5 ppm . Experimental values are for liquid-phase molecules unless otherwise stated.

|  | $\delta\left({ }^{15} \mathrm{~N}\right)($ calc. $)$ | $\delta\left({ }^{15} \mathrm{~N}\right)($ exp. $)$ |
| :--- | :---: | :---: |
| $\mathrm{NH}_{3}$ | 0.0 | 0.0 |
| HCN | 306.0 | 290.5 (gas) |
| $\mathrm{N}_{2}$ | 346.6 | 364.5 (gas) |
| NP | 681.4 | 612.5 (gas) |
| $\mathrm{NSF}^{2}$ | 720.5 | 264.7 |
| $\mathrm{NF}_{3}$ |  | 437.4 |
| $\mathrm{~N}_{2} \mathrm{~F}_{2}$ | trans | 506.7 |
| $\mathrm{H}_{2} \mathrm{CN}_{2}$ | $\mathrm{~N}(\mathrm{C})$ | 314.0 |
|  | N | 441.3 |

Table 7d Relative shifts $\delta\left({ }^{17} \mathrm{O}\right)$ referenced to $\mathrm{H}_{2} \mathrm{O}$ (liq.); $\mathrm{H}_{2} \mathrm{O}$ (gas) was used as primary reference for the calculated data, a value of $\delta\left(\mathrm{H}_{2} \mathrm{O}\right)=-36.1 \mathrm{ppm}$ was used for conversion, the calculated shielding $\sigma\left({ }^{17} \mathrm{O}\right)$ of $\mathrm{H}_{2} \mathrm{O}$ is 329.8 ppm.

|  | $\delta\left({ }^{17} \mathrm{O}\right)($ calc. $)$ | $\delta\left({ }^{17} \mathrm{O}\right)$ (exp.) |
| :--- | :---: | :---: |
| $\mathrm{H}_{2} \mathrm{O}$ | -36.1 | -36.1 (gas) |
| $\mathrm{H}_{3} \mathrm{O}^{+}$ | -10.3 | 9.1 (liq.) |
| $\mathrm{CH}_{3} \mathrm{OH}$ | -30.8 | -38 |
| CO | 368.1 | 350 (gas) |
| $\mathrm{CO}_{2}$ | 78.5 | 64.5 (gas) |
| $\mathrm{COS}^{-}$ | 218.2 | 200 (gas) |
| $\mathrm{NO}_{3}{ }^{-}\left(\mathrm{C}_{3}\right)$ | 437.3 | 410 (liq.) |
| $\mathrm{OSCl}_{2}$ | 374.6 | 292 (liq.) |
| $\mathrm{H}_{2} \mathrm{O}_{2}$ |  | 189.0 |
| $\mathrm{O}_{3}$ | O (cent.) | 1375.5 |
| $\mathrm{SO}_{2}$ | O (term.) | 1835.2 |

Table 7e Relative shifts $\delta\left({ }^{19} \mathrm{~F}\right)$ referenced to $\mathrm{CFCl}_{3}$; the calculated shielding $\sigma\left({ }^{19} \mathrm{~F}\right)$ of $\mathrm{CFCl}_{3}$ is 162.1 ppm . Experimental values are for gas-phase molecules unless otherwise stated.

|  | $\delta\left({ }^{19} \mathrm{~F}\right)($ calc. $)$ | $\delta\left({ }^{19} \mathrm{~F}\right)($ exp. $)$ |
| :--- | :---: | :---: |
| CFCl | 3 | 0.0 |
| HF | -247.6 | 0.0 |
| $\mathrm{BF}_{3}$ | -143.3 | -214.4 |
| $\mathrm{~F}_{2}$ | 359.6 | -131.6 |
| $\mathrm{CH}_{3} \mathrm{~F}$ | -300.9 | 428.4 |
| ClF | -501.1 | -275.4 |
| $\mathrm{SF}_{6}$ | 78.4 | -441.5 |
| $\mathrm{NF}_{3}$ | 139.2 | 56.1 |
| $\mathrm{PF}_{3}$ | -40.1 | 145.3 |
| $\mathrm{NSF}^{2}$ | 273.7 | -32.7 |
| $\mathrm{~N}_{2} \mathrm{~F}_{2}$ |  | 96.6 |

Table 7f Relative shifts $\delta\left({ }^{31} \mathrm{P}\right)$ referenced to $85 \% \mathrm{H}_{3} \mathrm{PO}_{4} ; \mathrm{PH}_{3}$ was used as primary reference for the calculated data, a value of $\delta\left(\mathrm{PH}_{3}\right)=-240 \mathrm{ppm}$ was used for conversion, the calculated shielding $\sigma\left({ }^{31} \mathrm{P}\right)$ of $\mathrm{PH}_{3}$ is 586.3 ppm .

|  | $\delta\left({ }^{31} \mathrm{P}\right)($ calc. $)$ | $\delta\left({ }^{31} \mathrm{P}\right)($ exp. $)$ |
| :--- | :---: | :---: |
| $\mathrm{PH}_{3}$ | -240 | -240 |
| $\mathrm{PH}_{4}{ }^{+}$ | -97.8 | -101 |
| PN | 325.1 | 302 |
| FCP | -321.2 | -207 |
| HCP | -9.0 | -32 |
| $\mathrm{PF}_{3}$ | 120.4 | 97 |
| $\mathrm{P}_{4}$ | -505.5 | -553 (sol.) |

Table $7 \mathbf{g}$ Relative shifts $\delta\left({ }^{33} \mathrm{~S}\right)$ referenced to $\mathrm{CS}_{2}$; the calculated shielding $\sigma\left({ }^{33} \mathrm{~S}\right)$ of $\mathrm{CS}_{2}$ is 526.4 ppm .

|  | $\delta\left({ }^{33} \mathrm{~S}\right)($ calc. $)$ | $\delta\left({ }^{33} \mathrm{~S}\right)$ (exp.) |
| :--- | :---: | :---: |
| $\mathrm{CS}_{2}$ | 0.0 | 0.0 |
| $\mathrm{H}_{2} \mathrm{~S}$ | -188.4 | -171.0 |
| $\mathrm{CH}_{3} \mathrm{SH}$ | -103.8 | -125 |
| COS | -257.9 | -261.8 |
| $\mathrm{SO}_{2}$ | 748.8 | 706.9 |
| $\mathrm{OSCl}_{2}\left(\mathrm{C}_{\mathrm{s}}\right)$ | 638.8 | 556 |
| $\mathrm{SF}_{6}$ | 212.8 | 155.4 (gas) |

Table 7h Several attempts for the calculation of $\delta\left({ }^{17} \mathrm{O}\right)$ relative shifts of the ozone molecule, O1-O2-O.

| Method | $\delta\left({ }^{17} \mathrm{O} 1\right)$ | $\delta\left({ }^{17} \mathrm{O} 2\right)$ |
| :--- | :---: | :---: |
| RHF/IGLO(II) | 3321 | 3216 |
| RHF/GIAO/6-311+G(2d) | 3624 | 3513 |
| MP2/GIAO/6-311+G(2d) | -2084 | -4246 |
| B3LYP/GIAO/cc-pvqz | 1835 | 1375 |
| Exp. | 1598 | 1032 |

the RHF wave functions. Therefore, ab initio MP2/6-311+G(2d) calculations were performed. However, entirely wrong results were obtained for $\mathrm{O}_{3}$. Finally, GIAO/B3LYP/cc-pvqz calculations yielded qualitatively correct results for $\mathrm{O}_{3}$, but the deviations from experimental data are still unusually high ( $200-300 \mathrm{ppm}$ ). A further source for remarkable deviations of $\delta\left({ }^{17} \mathrm{O}\right)$ data can be seen in the values from the liquid phase, for example, $\mathrm{SO}_{2}$ and $\mathrm{H}_{3} \mathrm{O}^{+}$. A difficult phosphorus-containing molecule is FCP, where RHF ab initio as well as DFT calculations fail by more than 100 ppm for the relative shift $\delta\left({ }^{31} \mathrm{P}\right)$.

Statistics for the deviations between calculated and experimental NMR chemical shifts $\delta$ are meaningful only under special conditions: 1) The number $N$ of investigated systems should be as large as possible. 2) The mean absolute deviation $|\Delta \delta|$ should be seen only in connection with the width of the global range of $\delta$.3) The breakdown of the calculation for some examples makes any statistics for the respective nucleus meaningless.

The mean absolute deviation for $\delta\left({ }^{1} \mathrm{H}\right)$ values in Table 7 a is as low as 0.4 ppm , but the range of $\delta\left({ }^{1} \mathrm{H}\right)$ is also very low ( 12.3 ppm ). The mean absolute deviation for $\delta\left({ }^{13} \mathrm{C}\right)$ values in Table 7 b is 9.2 ppm, where the range of $\delta\left({ }^{13} \mathrm{C}\right)$ is 342.7 ppm . For ${ }^{15} \mathrm{~N}$ in Table 7c, the disaster with the NSF molecule makes statistics meaningless. Statistics from the literature where NSF is not included are, therefore, misleading [17]. In addition, cations and anions are excluded from these statistics so that the respective $\delta$ ranges and, consequently, the mean absolute deviations are low, simulating best agreement between calculations and experiments. A similar situation can be found in Table 7 d for ${ }^{17} \mathrm{O}$, where the $\mathrm{O}_{3} \mathrm{~mol}$ ecule destroys any usefulness of statistics (cf. also Table 7h). For the nuclei ${ }^{19} \mathrm{~F},{ }^{31} \mathrm{P}$, and ${ }^{33} \mathrm{~S}$, the mean absolute deviations and, in parentheses, the widths of the $\delta$ ranges are (in ppm) 28.1 (869.9), 39.1 (855), and 37.3 (968.7), respectively. For example, the global $\delta$ range of ${ }^{31} \mathrm{P}$ for all presently known phosphorus compounds is as large as 2000 ppm . The utility of calculated NMR chemical shifts is restricted mainly to the description of trends of related compounds.

## 11. CONCLUDING REMARKS ON CALCULATED VERSUS EXPERIMENTAL PROPERTIES AND COMPUTATIONALLY DEMANDING MOLECULES

The meaning of this project can be demonstrated best by means of the pair of SiF and SiCl radicals. The four steps in chronological order in Table 8 are as follows: 1) The experimental dissociation energies were originally found in well-established tables [1,2].2) During the preparation of this project these radicals turned out to be bad $D_{0}$ systems in the sense of large discrepancies ( $\sim 0.5 \mathrm{eV}$ ) between experimental and B3LYP calculated $D_{0}$ values. 3) Members from the family of bad $D_{0}$ systems were selected for accurate time-consuming ab initio $\operatorname{CCSD}(\mathrm{T})$ calculations. As a result, these calculations favor the B3LYP calculated $D_{0}$ values over the experimental ones [88]. 4) Additional experimental $D_{0}$ data were found which are in excellent agreement with the B3LYP calculated values of this project. However, for SiF, these are not yet included in well-known tables such as [19] so that primary references [57] were late in finding, and for SiCl , these have been included in well-known tables [56] very recently, but primary references were not found.

The noble gas compound $\mathrm{He}_{2}{ }^{+}$exhibits large deviations between B3LYP calculated and accurate experimental properties. $A b$ initio $\operatorname{CCSD}(\mathrm{T})$ calculations confirm in this case the complete breakdown of the B3LYP procedure in its present form. Fortunately, $\mathrm{He}_{2}{ }^{+}$is unique in this respect. This is an important result from this study of a large number of molecules.

The bond length of the anion $\mathrm{Cl}_{2}{ }^{-}$is considered not to be well established from the experimental point of view. The B3LYP calculated bond length differs markedly from the experimentally estimated value. Recently performed $\operatorname{CCSD}(\mathrm{T})$ calculations did not solve the problem. Instead, a third quite dif-

Table 8 Properties of selected diatomic molecules in their chronological order.

| Molecule |  | $r_{\text {e }} / \AA$ | $\omega_{\mathrm{e}} / \mathrm{cm}^{-1}$ | $D_{0} / \mathrm{eV}$ | Source |
| :---: | :---: | :---: | :---: | :---: | :---: |
| SiF |  | 1.6011 | 857 | 5.57 | Exp. [1] |
|  |  | 1.6196 | 833 | 6.07 | B3LYP |
|  |  | 1.5997 | 863 | 6.14 | $\operatorname{CCSD}(\mathrm{T})$ [88] |
|  |  |  |  | 5.98 | Exp. [57] |
| SiCl |  | 2.058 | 536 | 3.85 | Exp. [1,2] |
|  |  | 2.0880 | 507 | 4.30 | B3LYP |
|  |  | 2.0568 | 535 | 4.47 | $\operatorname{CCSD}(\mathrm{T})$ [88] |
|  |  |  |  | 4.42 | Exp. [56] |
| $\mathrm{He}_{2}{ }^{+}$ |  | 1.0808 | 1698 | 2.36 | Exp. [1] |
|  |  | 1.1453 | 1360 | 3.27 | B3LYP |
|  |  | 1.0805 | 1701 | 2.35 | $\operatorname{CCSD}(\mathrm{T})$ |
| $\mathrm{Cl}_{2}{ }^{-}$ |  | $(2.65)^{\text {a }}$ | (249) ${ }^{\text {a }}$ | 1.26 | Exp. [1,35] |
|  |  | 2.7167 | 199 | 1.58 | B3LYP |
|  |  | 2.5526 | 271 | 1.29 | $\operatorname{CCSD}(\mathrm{T})$ [88] |
| MgO | ${ }^{3} \Pi$ | 1.870 | 648 |  | Exp. [44] |
|  | $\mathrm{X}^{1} \Sigma^{+}$ | 1.7490 | 785 | $(3.53)^{\text {a }}$ | Exp. [1] |
|  | ${ }^{1} \Sigma^{+}$ | 1.7389 | 821 |  | B3LYP |
|  | $\mathrm{X}^{3} \Pi$ | 1.8882 | 626 | 2.37 | " |
|  | $\mathrm{X}^{1} \Sigma^{+}$ | 1.7508 | 821 | $4.66{ }^{\text {b }}$ | ACPF [88] |
|  |  |  |  | $2.70{ }^{\text {c }}$ | " |
|  | $\mathrm{X}^{1} \Sigma^{+}$ |  |  | $2.56 \pm 0.2$ | Exp. [44] |

[^1]ferent value for the bond length is suggested [88]. Although we generally tend to trust $\operatorname{CCSD}(\mathrm{T})$ calculations, $\mathrm{Cl}_{2}{ }^{-}$could be one of the rare demanding systems for this approach as is chloryl chloride, $\mathrm{ClClO}_{2}$. This compound was investigated by rotational spectroscopy where the $\mathrm{Cl}-\mathrm{Cl}$ bond length $\left(r_{\mathrm{e}}\right)$ turned out to be $2.1921 \AA$ [89]. B3LYP calculations yielded a considerably larger value of $2.244 \AA$, but ab initio $\operatorname{CCSD}(\mathrm{T})$ calculations failed dramatically with $2.279 \AA$ [89].

Another difficult molecule is MgO. Experiments propose a ${ }^{1} \Sigma^{+}$ground state with an uncertain dissociation energy of 3.53 eV [1]. B3LYP calculations yield a ${ }^{3} \Pi$ ground state with a dissociation energy of only 2.37 eV . The deviation of 1.16 eV cannot be explained by the reversed electronic states because their relative energies differ by only 0.46 eV . ACPF ab initio calculations revealed two dissociation energies (cf. footnotes in Table 8). Later, an experimental thermodynamic dissociation energy (all atoms in their ground states) was found in the literature [44], which is in accordance with B3LYP as well as the ACPF calculations. Thus, the MgO problem is solved.

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[^0]:    ${ }^{*} 1 \mathrm{kcal} \xlongequal{ }$. 4.184 kJ

[^1]:    ${ }^{\mathrm{a}}$ Uncertain experimental values
    ${ }^{\mathrm{b}}$ Spectroscopic dissociation into ${ }^{1} \mathrm{~S}(\mathrm{Mg})+{ }^{1} \mathrm{D}(\mathrm{O})$
    ${ }^{\mathrm{c}}$ Thermodynamic dissociation into ${ }^{1} \mathrm{~S}(\mathrm{Mg})+{ }^{3} \mathrm{P}_{2}(\mathrm{O})$

